

1 **Beyond equilibrium: Re-evaluating physical modelling of fluvial systems to**
2 **represent climate changes**

3 Edwin R.C. Baynes^{1,*}, Wietse I. van de Lageweg^{1,+}, Stuart J. McLelland¹, Daniel R.
4 Parsons¹, Jochen Aberle^{2, 7}, Jasper Dijkstra³, Pierre-Yves Henry², Stephen P. Rice⁴,
5 Moritz Thom⁵, Frederic Moulin⁶

6

7 ¹ Geography and Geology, School of Environmental Sciences, Faculty of Science and Engineering,
8 University of Hull, Hull, HU6 7RX, UK.

9 ² Department of Civil and Environmental Engineering, Norwegian University of Science and
10 Technology, Trondheim, Norway.

11 ³ Marine and Coastal Systems, Department Ecology and Sediment Dynamics, Deltares,
12 Boussinesqweg 1, Delft, the Netherlands.

13 ⁴ Department of Geography, Loughborough University, Loughborough, LE11 3TU, UK.

14 ⁵ Forschungszentrum Küste, Leibniz Universität Hannover/TU Braunschweig, Hannover, Germany.

15 ⁶ Institut de Mecanique des Fluides de Toulouse (IMFT), Université de Toulouse, CNRS-
16 INPT-UPS, Toulouse, France.

17 ⁷ Leichtweiß Institute for Hydraulic Engineering and Water Resources, Technische Universität
18 Braunschweig, Braunschweig, Germany

19

20 * Now at: Université de Rennes, CNRS, Géosciences Rennes - UMR 6118, F-35000 Rennes, France

21 + Now at: Antea Group Belgium, Rodelveldlaan 1, 2600 Antwerpen, Belgium

22

1

23 **Abstract**

24 The interactions between water, sediment and biology in fluvial systems are complex and driven by
25 multiple forcing mechanisms across a range of spatial and temporal scales. In a changing climate,
26 some meteorological drivers are expected to become more extreme with, for example, more
27 prolonged droughts or more frequent flooding. Such environmental changes will potentially have
28 significant consequences for the human populations and ecosystems that are dependent on riverscapes,
29 but our understanding of fluvial system response to external drivers remains incomplete. As a
30 consequence, many of the predictions of the effects of climate change have a large uncertainty that
31 hampers effective management of fluvial environments. Amongst the array of methodological
32 approaches available to scientists and engineers charged with improving that understanding, is
33 physical modelling. Here, we review the role of physical modelling for understanding both biotic and
34 abiotic processes and their interactions in fluvial systems. The approaches currently employed for
35 scaling and representing fluvial processes in physical models are explored, from 1:1 experiments that
36 reproduce processes at real-time or time scales of 10^{-1} - 10^0 years, to analogue models that compress
37 spatial scales to simulate processes over time scales exceeding 10^2 - 10^3 years. An important gap in
38 existing capabilities identified in this study is the representation of fluvial systems over time scales
39 relevant for managing the immediate impacts of global climatic change; $10^1 - 10^2$ years, the
40 representation of variable forcing (e.g. storms), and the representation of biological processes.
41 Research to fill this knowledge gap is proposed, including examples of how the time scale of study
42 in directly scaled models could be extended and the time scale of landscape models could be
43 compressed in the future, the use of lightweight sediments, and innovative approaches for
44 representing vegetation and biostabilisation in fluvial environments at condensed time scales, such as
45 small-scale vegetation, plastic plants and polymers. It is argued that by improving physical modelling
46 capabilities and coupling physical and numerical models, it should be possible to improve
47 understanding of the complex interactions and processes induced by variable forcing within fluvial

48 systems over a broader range of time scales. This will enable policymakers and environmental
49 managers to help reduce and mitigate the risks associated with the impacts of climate change in rivers.

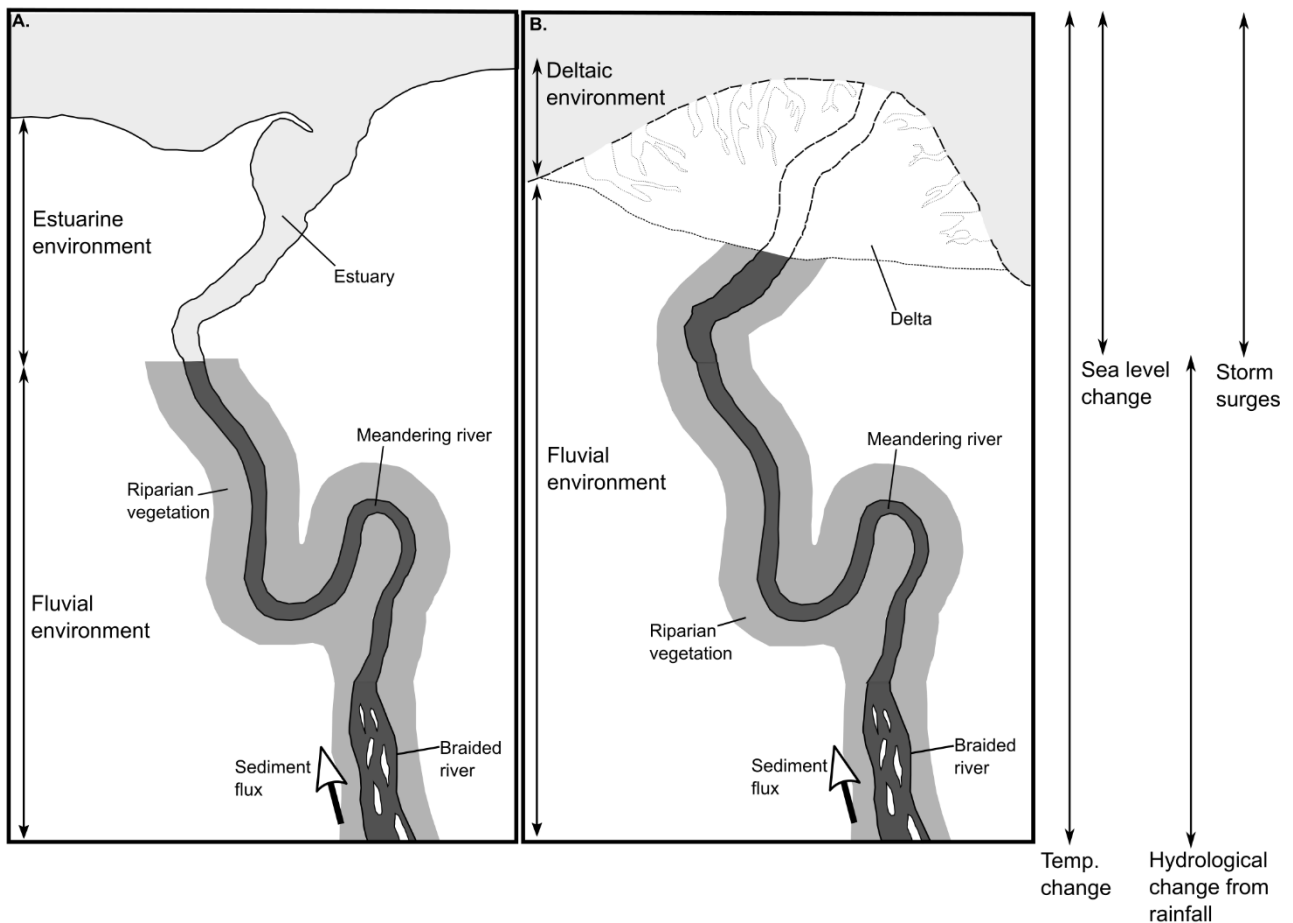
50

51 *Keywords: fluvial, climate change, physical modelling, scaling, review, floods, ecosystems, tipping*
52 *point*

53

54 1. Introduction

55 Global climate change is a grand challenge facing the Earth across numerous spatial and temporal
56 scales (IPCC 2014; EEA 2017) and the supply of water through the river networks is critically
57 important for the Earth's population (de Wit and Stankiewicz 2006). Expected impacts of climate
58 change in fluvial and fluvially-affected systems such as river deltas and estuaries (Figure 1) include
59 altered hydrological regimes and sediment fluxes (Nijssen *et al.* 2001; Syvitski *et al.* 2005), variations
60 in biota distribution and growth patterns (Harley *et al.* 2006), and more frequent extreme events such
61 as storm surges (Lowe & Gregory 2005), river floods (Garssen *et al.* 2015) and droughts (Garssen *et*
62 *al.* 2014). Understanding and adapting to these potentially irreversible and detrimental impacts
63 associated with new rates of environmental change and shifts in the frequency and magnitude of
64 events associated with climate change is therefore a fundamental priority for potentially vulnerable
65 fluvial environments, especially in regions where the human population is dependent on the local
66 water supply (de Wit and Stankiewicz 2006). In fact, management of fluvial environments presents
67 challenges in a changing climate, and requires an improved understanding of the feedbacks and
68 interactions between the driving mechanisms at work.



70

71 *Figure 1: Schematic diagram to highlight the environments within the scope of this review paper, with an*
 72 *estuarine environment shown in (A) and a deltaic environment shown in (B). Potential climate change*
 73 *impacts in these systems are identified. See Table 1 for details of expected changes in the environments*
 74 *induced by climate change.*

75 Physical modelling is an important tool in research on fluvial systems and an established
 76 technique for the design and testing of hydraulic structures. The high degree of experimental
 77 control in physical scale models allows for the simulation of varied, or rare, environmental
 78 conditions and hence measurements of conditions which cannot be measured in the
 79 prototype (i.e. the real site to be modelled). Moreover, physical modelling provides an
 80 essential link between field observations and theoretical, stochastic and numerical models
 81 which are required to predict the impact of environmental changes on aquatic ecosystems
 82 (Thomas *et al.* 2014). Physical modelling can therefore play a key role in the development
 83 of a better understanding of climate change impacts by improving our ability to predict these

84 impacts and, in turn, help adaptation to climate change-related challenges (Frostick *et al.*
 85 2011; Frostick *et al.* 2014).

86 *Table 1: Details of expected climate change induced impacts on fluvial and fluvial-affected estuarine and*
 87 *deltaic environments. Physical modelling studies can be used to understand these processes and test*
 88 *possible adaptation strategies.*

Climate change in forcing	induced mean	Predicted change	Associated impact on estuarine and fluvial environments	Source
Global surface temperature		By 2100: 0.3-1.7 °C temp rise (scenario RCP2.6*) 2.6-4.8 °C temp rise (scenario RCP8.5*)	Implications for vegetation growth in all environments	IPCC (2014)
Sea level rise		By 2100: 0.26-0.55 m (scenario RCP2.6) 0.45-0.82 m (scenario RCP8.5) 70% of coastlines worldwide experience change within 20% of global mean	Drowning of estuarine environments. Encroachment of saline water and associated impacts on biota. Increased aggradation of river deltas, accelerated channel and floodplain deposition and to higher channel avulsion frequency	IPCC (2014), Jerolmack (2009)
Storm surges		Largest increase in 50 year return period storm-surge height at UK coastline = 1.2 m (Scenario A2)	Increased risk from hazards (e.g. coastal flooding, coastal erosion) associated with storm surge events	Lowe and Gregory (2005)
Precipitation		Scenario RCP8.5: Increase in mean precipitation in high latitudes and equatorial Pacific. Decrease in mean precipitation in mid-latitudes Increase in extreme precipitation over most of mid-latitude landmasses and wet tropical regions become more intense and more frequent	Rivers: increased frequency and magnitude of higher peak flows, and possible prolonged drought periods with associated impacts for riparian vegetation distribution. Potential shifts in timing of seasonal hydrological regimes.	IPCC (2014), (Garssen <i>et al.</i> 2014; Garssen <i>et al.</i> 2015)
Waves		Latitude dependent: 0.6-1 m increase in 20 year return period wave height between 1990-2080 in NE Atlantic. Wave with 20 year return period in 1990 will have 4-12 year return period in 2080	Modification of the dynamics of estuarine and coastal systems	Wang <i>et al.</i> , (2004)

89 * RCP2.6 and RCP8.5 refer to two end-member Representative Concentration Pathways for anthropogenic
 90 greenhouse gas emissions. RCP2.6 refers to a stringent mitigation scenario, and RCP8.5 refers to a scenario
 91 with very high greenhouse gas emissions (IPCC, 2014).

92

93 Physical scale models are a key tool to simulate and investigate complex processes and
94 feedback mechanisms, with experimental designs that reflect the spatial and temporal scale
95 of the problem under investigation. Such techniques have been used for more than 100
96 years to investigate the interaction among flow, sediment transport, morphology, and
97 interactions with biota, enhancing the understanding of many different and complex
98 sediment transport and morphological processes across different spatial and temporal
99 scales (Kleinhans *et al.* 2015).

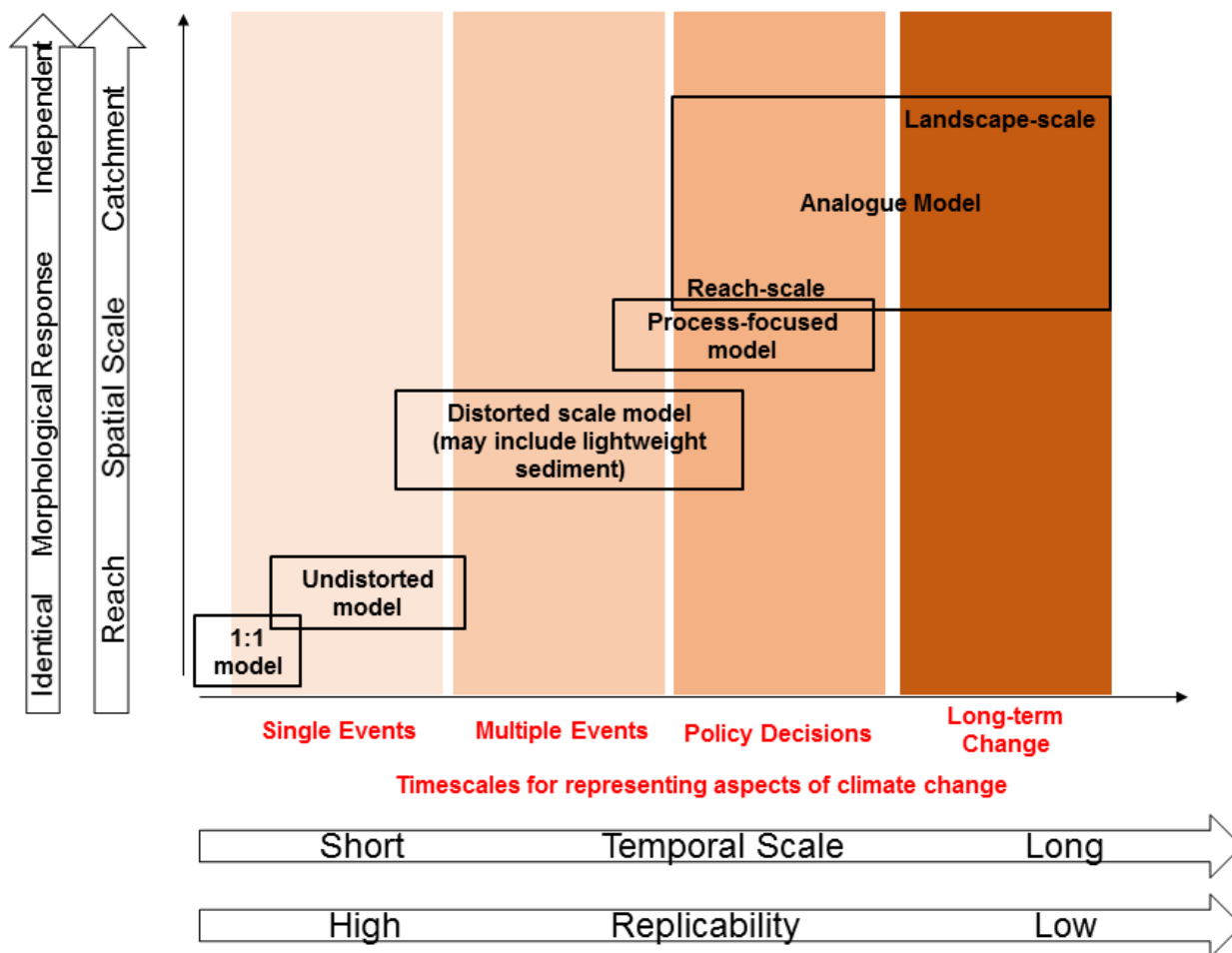
100 Physical modelling for climate change adaptation faces the challenge of incorporating, and
101 scaling, non-linear responses across a range of temporal and spatial scales resulting from
102 long-term changes in event frequency and magnitude. Until today, most physical models of
103 the impact of climate change on the aquatic environment have been based on stationary
104 boundary conditions reflecting a possible future climate state, with a simplified
105 representation of the systems (i.e. single grain size sediment, or no biotic elements).
106 Focusing solely on the final stage of a future scenario neglects that climate change may be
107 a progressive process which develops over time (IPCC 2014; EEA 2017). In particular, the
108 morphology of riverine, deltaic and estuarine environments will develop and change over
109 time in response to long-term changing boundary conditions. To address the challenges
110 related to climate change, it is crucial to develop an understanding of the complexity of the
111 systems, and how the environments adapt over longer periods of time, whether this change
112 is gradual or sudden, and how they behave under a different climate regime.

113 This review will examine current techniques and capabilities in physical modelling experiments for
114 representing climate change induced impacts on aspects of fluvial systems such as hydrodynamics,
115 sediment transport and morphodynamics. Firstly, this review provides a technical discussion of
116 different modelling approaches and the formal scaling laws that they obey (section 2), before
117 identifying the challenges that physical models face for representing variable forcing and the impacts
118 of climate change within experiments (section 3). Section 4 provides detailed examples of recent

119 innovative approaches at the forefront of the physical modelling of environmental systems and how
120 these facilities might be developed in the future.

121 **2. Scaling approaches and challenges in representing different time** 122 **scales in physical modelling**

123 Figure 2 presents a schematic overview of different model types and their ability to replicate
124 the relevant spatial and temporal scales of the prototype. In the discussion below, we explain
125 the essence of each of these approaches, the scaling laws that they must successfully
126 achieve and provide some examples of their application for the understanding of fluvial
127 processes and systems.



128

129 *Figure 2. The relative application of different approaches for physical modelling, with different approaches*
 130 *being more appropriate for modelling processes over different spatial and temporal scales. Developed*
 131 *from Peakall et al. (1996).*

132 In scaled models, the time passes generally faster than in the prototype, which makes them
 133 attractive for the study of climate change impacts. However, as will be outlined below, their
 134 design and the interpretation of results can be challenging because the hydrodynamic time
 135 scales are generally quite different from those for morphodynamic fluvial adjustments
 136 (Tsujiimoto 1990), and the scaling of biota is even more uncertain. Models based on both
 137 geometrical and dynamic similarity (i.e. by scaling important force ratios; see below) are a
 138 well-established approach for designing hydraulic structures at larger spatio-temporal scales
 139 while distorted models (models with different geometrical scale ratios in the horizontal and

140 vertical directions), and relaxed-scale analogue models attempt to reproduce some selected
141 properties of the prototype (Peakall *et al.* 1996).

142 The scaling laws used to design physical models can be derived based on a dimensional
143 analysis (Buckingham 1914; Barenblatt 2003). An important prerequisite for the design of a
144 physical model is the dynamic similarity that ensures a constant prototype-to-model ratio of
145 the masses and forces acting on the system (Einstein and Chien 1956; Yalin and Kamphuis
146 1971; Hughes 1993; Frostick *et al.* 2011), i.e. that the derived dimensionless parameters
147 are equal in model and prototype. Important force ratios defining these dimensionless
148 numbers can be obtained by considering inertia, gravity, viscosity, surface tension, elasticity
149 and pressure forces, respectively. A perfect dynamic similarity for all possible force ratios
150 cannot normally be achieved for model scales that deviate from the prototype scale since
151 the same fluid (water) is used in both prototypes and models. This means that it is not
152 possible to design a downscaled model so that the relative influence of each individual force
153 acting on a system remains in proportion between prototype and model as outlined by e.g.,
154 de Vries, (1993); de Vries *et al.* (1990); Ettema & Muste (2004); Heller (2011); Hudson
155 (1979); Hughes (1993); Sutherland & Whitehouse (1998) and Yalin (1971). Scale models
156 need therefore to be designed in a way that maintains important force ratios whilst providing
157 justification for neglecting other force ratios. Neglecting force ratios will result in scale effects
158 if the model is operated at boundary conditions where the neglected force ratios are
159 important; in other words, there will be a divergence between up-scaled model
160 measurements and real-world observations. Scale effects become more significant with
161 increasing scale ratio and their relative importance depends on the investigated
162 phenomenon (Heller 2011), i.e. scale effects will have to be accepted.

163 In the following discussion of the different modelling approaches, it is assumed that the
164 model studies are carried out with water as model fluid so that the ratio of fluid properties in
165 model and prototype such as fluid density ρ_r , fluid dynamic and kinematic viscosity μ_r and ν_r ,

10

166 respectively are equal to 1; the subscript r denotes the ratio between model (m) and
167 prototype (p). Moreover, scale effects due to fluid temperature will not be considered
168 although it is worth mentioning that Young and Davies (1991) used heated water (30° C) in
169 their experiments in order to achieve closer similarity in Reynolds numbers. Finally, although
170 beyond the scope of this review, experiments using dense fluids have been used to study
171 grains at the threshold of motion (e.g. oil (Best 1998) and glycerol (Guerit *et al.* 2014)), and
172 scaling for morphological processes in extra-terrestrial environments is also possible (e.g.,
173 aeolian dunes on Mars and Venus (Claudin and Andreotti 2006); morphological
174 development on Mars (Kraal *et al.* 2008; Dietrich *et al.* 2017).

175 **2.1.1:1 Physical models**

176 Models that replicate the prototype with no reduction in dimensions can be described as 1:1
177 physical models (Figure 2). 1:1 models are mainly used to study physical processes at the
178 smallest spatial and temporal scales under controlled conditions. Examples include
179 experiments aimed at replicating flow turbulence structures in open channels to predict
180 incipient motion and sediment transport (Shields 1936; Zanke 2003; Nikora *et al.* 2001;
181 Hofland, Battjes, and Booij 2005; Grass 1971). Full-scale replication of the larger
182 components of rivers such as channels, levees and bars requires a lot of space with
183 associated high operational costs and these experiments are therefore rare. An example of
184 a 1:1 model is provided by the Smart Levee project in which a river dike is replicated (Figure
185 3, <http://www.floodcontrolijkdijk.nl/en/experiments>). The full-scale physical model allows for
186 experiments on piping, micro- and macro-stability, and flow slide in the absence of scale
187 effects.

188



189

190 Figure 3. Example of a physical model of a river dike taking a 1:1 approach
 191 (<http://www.dijkmonitoring.nl/en/projects/>).

192 **2.2. Undistorted models**

193 Geometrical similarity means that all scales with dimensions of length x , y , z are equal: $r_x =$
 194 $r_y = r_z$ (Figure 2), and in *undistorted models* the geometry of the model is consistent with the geometry
 195 of the prototype. The most commonly used scaling approach for fluid flow in *undistorted*
 196 *models* is Froude-scaling, which requires similarity in the Froude number in model and
 197 prototype:

198
$$Fr = U/(gh)^{0.5} \tag{1}$$

199 where U denotes the mean flow velocity, g the gravitational acceleration, and h the water
 200 depth. This scaling law, ensuring the constant ratio between inertia and gravitational forces
 201 in model and prototype, is most significant for open channel flows and ensures that the water
 202 surface will be adequately replicated in the model (Kobus 1978).

203 Considering a uniform open-channel flow with a fixed bed in a wide channel, i.e. a width to
204 depth ratio > 30 so that the hydraulic radius can be replaced by the water depth h ,
205 dimensional analysis results in four important dimensionless parameters, which are the
206 Froude number as defined above, the flow Reynolds number $Re = Uh/\nu$, the relative
207 roughness k/h (with k = roughness length scale, which is often expressed in terms of the
208 grain diameter d), and the slope S . Requiring Froude number similarity in model and
209 prototype means that the flow Reynolds-number (Re) will differ between the model and
210 prototype (it can be shown that for Froude-scaled models $Re_m < Re_p$). To avoid
211 corresponding scale effects, the flow in both model and prototype needs to be fully turbulent
212 so that viscosity effects are negligible. The roughness (or friction losses) can be scaled
213 considering similarity in the Darcy-Weisbach friction factor or alternatively in the Chézy-
214 coefficient or Manning number, and the model slope equals by definition the prototype slope
215 in undistorted models.

216 Movable bed models represent a two-phase flow with a solid (particles) and fluid phase
217 (Yalin 1959). While the flow is generally Froude-scaled, the similarity in sediment movement
218 depends on a set of additional dimensionless parameters which are the grain Reynolds-
219 number $Re^* = v_*d/\nu$, densimetric Froude number (Shields-number) $Fr^* = \rho v_*^2 / [(\rho_s - \rho)gd]$,
220 relative sediment density ρ_s/ρ , relative submergence h/d , and relative fall speed v_s/v_* (see
221 Yalin 1971; Hughes 1993; Peakall *et al.* 1996 for details). Peakall *et al.* (1996; 2007) argued
222 that the 90th percentile of the sediment grain size (D_{90}) should be used in the calculation of
223 the grain Reynolds-number (Re^*), as the coarsest grains contribute the most to the definition
224 of the hydraulic conditions due to their impact on the roughness of the sediment surface.
225 Recently, Kleinhans *et al.* (2017) argued that percentiles lower than the D_{90} can be used
226 when the sediment mixture contains a wide range of grain sizes, as long as the percentile
227 used protrudes above the viscous flow sublayer to contribute to roughness. In these
228 definitions, ρ_s denotes the sediment density, v_* the shear velocity, and v_s the fall velocity.

229 To obtain perfect similitude for sediment transport processes in model studies using water
 230 as fluid (i.e., $\rho_r = \nu_r = 1$), all these quantities would have to be equal in the model and
 231 prototype resulting in:

$$232 \quad Re_{*r} = d_r v_{*r} = 1 \quad (2)$$

$$233 \quad Fr_{*r} = \frac{v_{*r}^2}{(\rho_s - \rho)_r d_r} = 1 \quad (3)$$

$$234 \quad \rho_{s,r} = 1 \quad (4)$$

$$235 \quad \frac{h_r}{d_r} = 1 \quad (5)$$

$$236 \quad \frac{v_{s,r}}{v_{*r}} = 1 = \frac{\sqrt{L_r}}{t_r} \quad (6)$$

237 Equations (2) – (6) were formulated for unidirectional flow conditions for which the shear
 238 velocity can be determined via $v_* = (ghS)^{0.5}$ so that, for example, equation (3) can be written
 239 as:

$$240 \quad Fr_{*r} = \frac{h_r^2}{(\rho_s - \rho)_r L_r d_r} = 1 \quad (7)$$

241 A general problem encountered in the scaling of shear velocity v_* (or bed shear stress) is
 242 that this similitude assumes a flat bed. This is not necessarily the case because the bed
 243 topography of most riverine environments is characterized by bedforms or other
 244 morphological features (Hughes 1993), i.e. scale effects may be induced if such
 245 morphological features are not adequately reproduced or if d_r deviates from the vertical
 246 scale ratio h_r (Gorrick and Rodríguez 2014). Based on the similarity in Fr it becomes possible
 247 to derive the hydraulic time scale t_r (Kobus 1978):

$$248 \quad t_r = \frac{L_r}{\sqrt{h_r}} \quad (8)$$

249 For a non-distorted model $t_r = L_r^{0.5}$, indicating that time related to bulk flow processes in the
 250 model passes faster than in the prototype.

251 The mechanism for suspended sediment transport differs from the mechanism for bed load
252 transport. This is reflected by the criterion defined by equation (6) corresponding to the ratio
253 of settling velocity to shear velocity, i.e. the Rouse number, which is most important for
254 suspension-dominated models. Such models are more common in coastal modelling
255 applications than in alluvial river studies and require the reproduction of the uplift of particles
256 due to turbulence induced by waves or currents, and their subsequent transport in the water
257 column. In this context it is worth mentioning that, in the case of waves, such models require
258 the consideration of different physical parameters in equations (2) - (6) than fluvial bed load
259 models, such as the characteristic velocity $(gH_b)^{0.5}$ instead of the shear velocity v^* and the
260 breaking wave height H_b instead of water depth h (Hughes 1993).

261 Assuming Froude-similarity for the flow and inserting the corresponding hydraulic time scale
262 given by equation (8) into equation (6) yields:

$$263 \quad h_r = L_r \quad (8)$$

264 i.e. the dynamics of the suspended load transport can only be modelled exactly using an
265 undistorted model. Considering all scaling criteria, it is therefore only possible for one
266 transport mode to be modelled following similarity criteria while the other mode will be
267 affected by scale-effects (Hughes 1993). Nonetheless, physical model experiments that
268 simulate both modes of sediment transport have been attempted (Grasso *et al.* 2009). If
269 movable bed models need to be distorted, the distortions should not be so large that the
270 type of sediment transport changes (i.e. from bed load to suspended load or *vice versa*).

271 When maintaining the similarity in sediment density ($\rho_{s,r} = 1$ or $(\rho_s - \rho)_r = 1$), *undistorted*
272 *models* fulfill the criteria given by eqns. (3) to (5) while violating the fall velocity (equation 6)
273 and the grain-Reynolds number criterion (equation 2). The latter corresponds for this model
274 type to $Re_{\tau} = L_r^{1.5}$ indicating that they should be operated in hydraulic rough conditions, i.e.
275 $Re^* > 70$, to avoid scale effects arising through viscous forces as Re^* in prototype conditions

276 will be larger than in the model. Recent work has indicated that the value of $Re^* > 70$ to
277 define hydraulically rough conditions may be overly conservative, with the value potentially
278 as low as 15 (Kleinhans et al., 2017). An important limitation of this type of model in regard
279 to the scale factor arises from the requirement to scale the sediment with the same factor
280 as the model length scale. If, for example, fine sand is already present in the field, then this
281 restriction could easily result in the requirement to use potentially cohesive sediments
282 representing a problem due to the different properties of cohesive sediments compared to
283 granular material. To minimize this problem, special materials may be used such as
284 Ballotini® (non-cohesive glass microspheres with diameters as small as 45 μm) or different
285 model types as described below.

286 **2.3. Distorted physical models**

287 Distorted models are characterised by different horizontal and vertical length scales so that
288 $S_r \neq 1$ (Figure 2). The distortion leads directly to scale effects in the flow field (see e.g. Lu *et*
289 *al.* 2013; Zhao *et al.* (2013) and geometric similarity may be replaced by geometric affinity
290 (De Vries 1993). Distortion is not acceptable in a model where the vertical velocity
291 components are important, but vertically distorted models are acceptable for uniform, non-
292 uniform and unsteady flow conditions with relatively slow vertical motion (Novak *et al.* 2010).
293 For example, considering scale models of rivers, the horizontal dimensions involved are
294 commonly much larger than the vertical dimensions and this will lead to unrealistic scale
295 models if the vertical scale ratio (h_r) is selected equal to the horizontal length scale ratio (L_r)
296 (De Vries 1993). Additional care needs to be taken in regard to potential scale effects due
297 to water surface tension if the water depth in the model is low (Hughes 1993; Peakall and
298 Warburton 1996; Van Rijn *et al.* 2011) or if the model is operated with varying background
299 water levels (e.g., to simulate tidal effects) because the effect of wetting and drying bank
300 material will change its behaviour (e.g, Thorne and Tovey, 1981).

301 The key issue in reproducing mobile bed morphology is sediment mobility. Particle size
302 cannot be reduced to the same degree as the other x, y, z dimensions of the experiment
303 relative to the prototype because properties such as incipient motion and cohesion of silt
304 and clay are significantly different from those of sand and gravel (Lick and Gailani 2004).
305 Given the small water depth and flow velocities in this model type, sediment mobility is
306 typically lower than in the prototype or may even be below the beginning of sediment motion.
307 Three methods have classically been applied to overcome this issue (Kleinhans *et al.* 2014):
308 i) a vertical distortion of the model leading to increased gradients and a decrease of surface-
309 tension effects (Peakall *et al.* 1996); ii) tilting of the bed, which further increases the gradient;
310 or iii) the introduction of lightweight sediment.

311 Vertical exaggeration of the model compared to the prototype has a range of effects on
312 sediment transport, morphodynamics and resultant stratigraphy. Stronger bed gradients
313 combined with small water depths affect the threshold for the beginning of sediment motion
314 (Shields 1936; Vollmer and Kleinhans 2007), which cascades into differences in sediment
315 sorting patterns between the model and the prototype (Solari and Parker 2000; Seal *et al.*
316 1997; Toro-Escobar *et al.* 2000; Wilcock 1993; Peakall *et al.* and Best 2007; Stefanon *et al.*
317 2010). In addition, it can be shown analytically that wavelengths, migration rates and
318 amplitudes of river bars are a function of channel width-to-depth, sediment mobility as well
319 as channel curvature, width variations and sinuosity (Struiksma 1985; Seminara and Tubino
320 1989; Talmon *et al.* 1995). This implies that any vertical distortion in the scale model will
321 alter the morphology and resultant stratigraphy as seen in the prototype.

322 The introduction of lightweight sediments results in similarity in both Re^* and Fr^* while
323 violating intentionally the sediment density as well as the relative roughness criterion. As
324 indicated by the name, this type of models makes use of model sediments with a lower
325 density than the prototype sediment. For models focusing on bed load transport it may be
326 reasonable to relax the criterion defined by equation (6). Low (1989) found in experiments
17

327 with lightweight materials of different specific densities $1 < \rho_s/\rho < 2.5$ and a grain diameter
328 of $d = 3.5$ mm that the specific volumetric bed load transport rate q_s was related to $v_{*r}/v_{s,r}$
329 by a simple power relation and that $q_s \sim v_{*r}^6$ and $\sim v_{s,r}^{-5}$. Zwamborn (1966) argued that the
330 Fr^* criterion (equation 3) is essentially the same as the $v_{*r}/v_{s,r}$ -criterion and that a good
331 similarity in river morphology can be expected between model and prototype if the latter
332 criterion is used together with an appropriate friction criterion and near similarity in Re^* . More
333 details in regard to the scaling laws considering or neglecting the fall-speed dependency for
334 such models can be found in Hughes (1993) and Van Rijn *et al.* (2011).

335

336 **2.4. Process-focused physical models**

337 Here we introduce the term process-focused physical models (Figure 2) to describe
338 *Densimetric Froude models* that relax the similitude in Re^* (Eq. 2) whilst maintaining
339 similarity in Fr^* (Eq. 3), but do not have a particular target natural prototype in mind. These
340 models allow the investigation of the processes and generic planform morphologies such as
341 channel braiding by reproducing fundamental sediment transport processes such as
342 bedload transport and exploring the sensitivity of processes and morphologies to different
343 experimental conditions. Bed sediment must be mobile in the bedload regime to replicate
344 gravel-bed rivers in nature and mobile in the suspension regime to replicate sandbed rivers,
345 which is challenging due to cohesive effects for silt and clay if used to represent scaled down
346 sand (Smith 1998; Hoyal and Sheets 2009). This class of models simplifies the
347 representation of both discharge regimes and sediment properties using simple flow regimes
348 (constant discharge or single events to represent annual floods) and a hydraulically rough
349 bed to minimise scale effects, which conflicts with sediment mobility requirements. This
350 conflict is generally solved by applying a poorly sorted sediment mixture in which the
351 coarsest fraction ensures hydraulic rough conditions (Peakall *et al.* 2007; Van Dijk *et al.*

352 2012). Examples of process-focused models include the experiments aimed at river
353 meandering by Friedkin (1945) and the braided river experiments by Ashmore (1988). Many
354 practical applications of such models indicate their suitability in studying morphodynamic
355 processes within river reaches as well as for coastal environments (Hughes 1993; Willson
356 *et al.* 2007; Kleinhans *et al.* 2014).

357 There is an overlap between distorted models and process-focused models when similitude
358 in Re^* may be close to specific natural prototype situations (Figure 2). Similarly, the point at
359 which a process-focused model should be described as an analogue physical model is not
360 always clear since it is not known when simplifications in sediment characteristics or
361 discharge regimes make model behaviour differ significantly from a natural system.

362

363 **2.5. Analogue physical models**

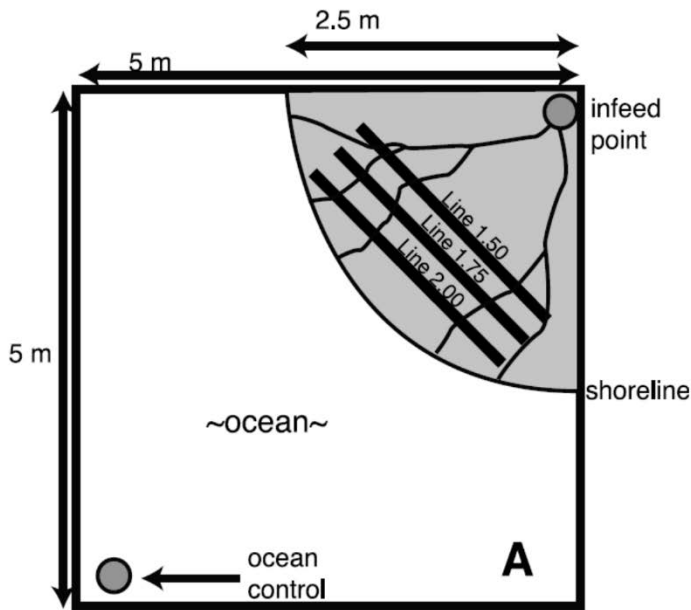
364 The evolution of river morphodynamics over larger spatial and temporal scales is often
365 investigated in so-called *analogue models* (Davinroy *et al.* 2012), which are designed to
366 represent larger prototype environments over longer periods of time (Figure 2). *Analogue*
367 *models* are designed to study analogies or ‘similarity of process’ between the model and
368 prototype and are not designed to keep strict similarity in the above scaling criteria (Hooke
369 1968), although they can theoretically be classified according to the model types defined
370 above. However, the aforementioned model types are generally stricter in terms of similarity
371 criteria than *analogue models* for which the validation or ‘effectiveness’ (Paola *et al.* 2009)
372 depends on the judgement of similitude in bed-sediment movement (Ettema & Muste 2004)
373 or on the operator due to the lack of a specific methodology for describing the degree of
374 morphodynamic and stratigraphic similarity in model studies (Gaines and Smith 2002). Yet,
375 well-designed analogue models have been shown to be an essential tool for studying
376 morphodynamic processes and stratigraphic expressions across a wide range of spatial

377 scales for different river channel morphologies and fluvially-affected coastal environments
378 (Bruun 1966; Hudson 1979; Peakall *et al.* 2007; Green 2014; Bennett *et al.* 2015; Yager *et*
379 *al.* 2015), despite violating the aforementioned scaling rules in many ways (Paola *et al.* 2009;
380 Kleinhans *et al.* 2014; Peakall *et al.* 1996; Kleinhans *et al.* 2015).

381 Due to the large range in spatial and temporal scales covered by *analogue models*, two sub-
382 groups can be identified (Figure 2). First, *analogue-reach* scale models are process-focused
383 physical models with an added degree of scaling relaxation. Examples include the
384 introduction of alfalfa as vegetation into the models as a representation of vegetation effects
385 in nature. A host of experiments has highlighted the important role vegetation can have in
386 controlling bank erosion and river pattern formation under the simplest conditions (Tal &
387 Paola 2007; Tal & Paola 2010; Braudrick *et al.* 2009; Van de Lageweg *et al.* 2010; van Dijk
388 *et al.* 2013a). The addition of fine silica flour in the experiments of Peakall *et al.* (2007) and
389 Van Dijk *et al.* (2013b) as the finest sediment into the models as a representation of cohesive
390 silt and clay in nature can also be considered an analogue-reach modelling approach, and
391 has been shown to lead to active meandering systems due to the added cohesion of
392 incorporating fine grained material (Peakall *et al.*, 1996; 2007; Kleinhans *et al.*, 2014). The
393 addition of nutshells has been used to represent low-density and highly-mobile sediment
394 acting as floodplain filler (Tambroni *et al.* 2005; Hoyal and Sheets 2009; Van de Lageweg
395 *et al.* 2016; Ganti *et al.* 2016). Similarly, a wide range of extracellular polymeric substances
396 (EPS) has been introduced into models to represent biological cohesion (Hoyal & Sheets
397 2009; Kleinhans *et al.* 2014; Schindler *et al.* 2015; Parsons *et al.* 2016). For example, EPS
398 has been used in analogue delta experiments to increase the range of natural
399 morphodynamics processes that can be reproduced, by increasing the cohesion of the
400 sediment material (Hoyal and Sheets, 2009). The polymer-sediment mix, developed at the
401 ExxonMobil Upstream Research Company (Hoyal and Sheets, 2009) performed best in the
402 presence of clay and sand, and the deltas produced during the experiments had geometries

403 characteristic of natural deltas composed of sandy noncohesive sediments, allowing
404 experimental investigations of forcing factors such as sea-level rise on channel mobility and
405 shoreline dynamics (Martin et al., 2009).

406 Second, *analogue-landscape* models represent the spectrum of scale models associated
407 with the largest spatial and temporal scales shown towards the top right in Figure 2. Such
408 models typically concern an entire landscape (e.g. delta or mountain range) and aim to
409 explore its evolution across longer (e.g. geological) time scales. River-delta landscape
410 experiments provide an example of this type of scale model (Figure 4). The analysis of this
411 experimental data allowed the identification of a small, but significant, chance for the
412 preservation of extreme events in the stratigraphy due to the heavy tailed statistics of
413 erosional and depositional events (Ganti *et al.* 2011). This quantified understanding of the
414 evolution of a river delta system under rising base level would only be possible using the
415 *analogue-landscape* modelling approach, where processes characteristic of larger delta
416 systems are replicated and monitored at high spatial and temporal resolutions that would be
417 impossible in the field.



418

419 Figure 4. Example of an experiment using an analogue-landscape modelling approach (Sheets,
 420 Hickson, and Paola 2002; Ganti *et al.* 2011). (a) Schematic of the experimental set up. (b)
 421 Photography of the delta after 11 hours of experimental run time. From Ganti *et al.* (2011).

422

423 3. Challenges representing climate change impacts in physical models

424 The impacts of climate change, and more broadly, non-constant forcing, will affect fluvial systems
 425 over a range of time scales. Increased magnitude of individual events to millennial-scale shifts in
 426 long-term forcing dynamics such as the total volume and seasonal variations in annual precipitation
 427 and changes in the biological characteristics could have dramatic impacts on the state and

428 functionality of fluvial systems (Wobus, Tucker, and Anderson 2010). This section identifies the
429 current challenges in representing these impacts on the fluvial environment using physical models.

430 **3.1. Differing timescales of morphodynamic and hydrodynamic processes**

431 Hydrodynamic processes usually occur at a much shorter time scale than morphodynamic
432 processes and, as will be shown below, time scales related to different morphological
433 processes do not necessarily coincide in physical models (Yalin 1971). This can, in turn,
434 result in undesired scale-effects that become more significant with decreasing physical
435 model scale (i.e. of the reproduction of the prototype) (Figure 2).

436 The determination of sedimentological time scales in movable-bed models is difficult and
437 often subjective. In fact, the sedimentological time scale cannot be freely chosen as it results
438 from the chosen scales of the other model parameters (Hentschel 2007) and, hence,
439 depends on which scaling criteria are intentionally violated. Moreover, there is the need to
440 distinguish between different time scales for different morphological processes such as
441 individual grain movement ($t_{sg,r}$) and the evolution of the bed surface in the vertical (t_{η}) and
442 horizontal (t_{Lr}) directions, respectively. Corresponding time scales are presented in general
443 terms in Table 2.

444

445 Table 2. Time scales for bed load dominated models, $\mu_r = \nu_r = 1$, and assuming $v^* = (ghS)^{0.5}$.

Time scale	Eq.	Criteria and comments	Source
$t_{sg,r} = d_r L_r^{0.5} h_r^{-1}$	(10)	- individual grain movement	Yalin (1971)
$t_{sg,r} = L_r h_r^{-2}$	(11)	-individual grain movement - similarity in Re^*	Yalin (1971)
$t_{\eta r} = L_r h_r$	(12)	- similarity in dimensionless transport rate -similarity in Re^* - porosity equal in model and prototype	Yalin (1971)
$t_{\eta r} = L_r^{1.5} d_r^{-1} (1-\phi)_r$	(13)	- similarity in dimensionless transport rate	Hentschel (2007)
$t_{\eta r} = L_r^{2.5} h_r^{-2} (1-\phi)_r (\rho_s - \rho)_r$	(14)	- similarity in dimensionless transport rate - similarity in Fr^*	Hentschel (2007))
$t_{\eta r} = L_r h_r^{1.5} d_r^{-7/6} (1-\phi)_r$	(15)	- similarity in dimensionless transport rate -similarity in Fr^* - near similarity in Re^*	Tsujimoto (1990)
$t_{Lr} = L_r^{1.5} h_r^{-1}$	(16)	- individual grain movement	Yalin (1971)

446

447 According to Yalin (1971), the movement of an individual bed load grain is governed by the
 448 geometrical scale of the particle diameter d and the kinematic scale v^* , respectively resulting
 449 in the time scales $t_{sg,r}$ defined by equations (10) and (11), where equation (12) results from
 450 the additional requirement of similarity in Re^* .

451 Considering the temporal development of a movable bed surface in a physical model,
 452 different scales in the horizontal and vertical directions need to be taken into account. For
 453 fluvial environments, the most common approach to derive the time scale for the formation
 454 of a movable bed surface is based on the comparison of the model response time to known
 455 prototype response times (Vollmers and Giese 1972; Kamphuis 1975; Einstein and Chien
 456 1956). This is typically achieved by considerations of the variation of the bed surface level
 457 η in vertical direction with time and the volumetric sediment transport rate q , i.e. the Exner
 458 equation (Paola and Voller 2005; Coleman and Nikora 2009). Thus, the corresponding time
 459 scale can be defined according to Tsujimoto (1990) and Hughes (1993):

$$460 \quad t_{\eta r} = \frac{L_r h_r (1 - \phi)_r}{q_r} \quad (17)$$

461 where ϕ denotes the porosity of the bed material. A similar formulation can be obtained
 462 considering the movement of river dunes assuming their geometrical similarity in model and
 463 prototype. Introducing the dimensionless volumetric bed load transport rate $q^* = q / (v^* d)$,
 464 equation (17) can be rewritten according to:

$$465 \quad t_{\eta r} = \frac{L_r h_r (1 - \phi)_r}{q^*_r d_r v^*_r} \quad (18)$$

466 Assuming similarity in q^* in model and prototype (i.e. $q^*_r = 1$), equation (18) represents the
 467 basis for equations (12) to (15) in Table 2 for which it was assumed that $v^*_r = (h_r S_r)^{0.5} = h_r L_r^{-0.5}$.
 468 Note that for geometrically similar grains with a similar grain-size distribution $(1 - \phi)_r = 1$
 469 (Hentschel 2007). Also, for practical purposes, the sediment transport rate is often
 470 determined from existing bed load formulae. Using such relationships in equation (18),
 471 instead of a measured q^* , can result in different time scale calculations.

472 Equation (16) in Table 2 was derived by Yalin (1971) and describes the time scale related
 473 to the evolution of the mobile bed surface in horizontal direction. This equation is based on
 474 single grain movement considerations and the relation of the diameter scale with the
 475 longitudinal scale.

476 Comparing the different time scales given in Table 2 it becomes apparent that

$$477 \quad t_{\eta r} < t_{Lr} < t_r < t_{sgr} \quad (19)$$

478 i.e. the vertical evolution of the bed surface has the shortest time scale, followed by the
 479 longitudinal displacement of the grains and the hydrodynamic time scale. The longest time
 480 scale is for the individual motion of a grain (Peakall, Ashworth, and Best 1996). Other time
 481 scales than those discussed here may be derived based on the consideration of the
 482 evolution of morphodynamic features such as meander bend migration rate, floodplain

483 evolution and biological development (Tal and Paola 2007; Kleinhans *et al.* 2014, and
484 references therein).

485 The time scales can also be linked to the bed-load models defined above. In *undistorted*
486 *similarity models* with unidirectional flow $t_{sg,r} = t_{\eta r} = t_{Lr} = L_r^{0.5}$, which is equal to the hydraulic
487 time scale t_r . *Geometric similarity models* therefore offer the opportunity to study the effects
488 of hydrographs on bed evolution. The time scales for *distorted lightweight models* can be
489 derived as $t_{sg,r} = (\rho_s - \rho)_r^{-2/3}$, $t_{\eta r} = h_r^3(1-\phi)_r (\rho_s - \rho)_r^{-2/3}$, $t_{Lr} = h_r^2(\rho_s - \rho)_r^{-1}$ thereby assuming
490 $q_{r^*} = 1$ and that bed shear stress can be determined from the depth slope product.

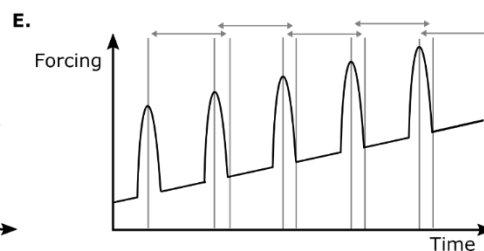
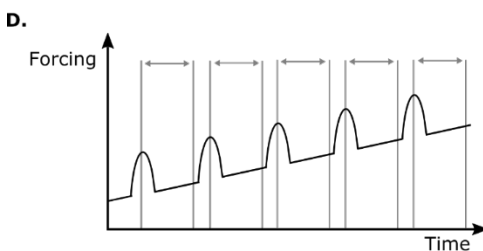
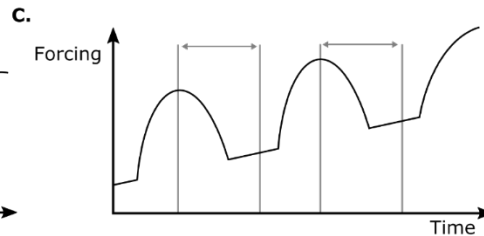
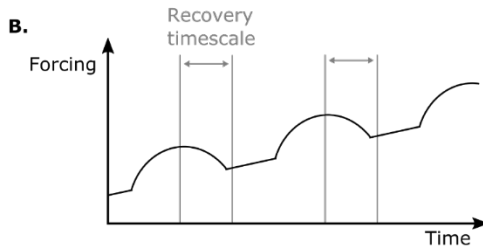
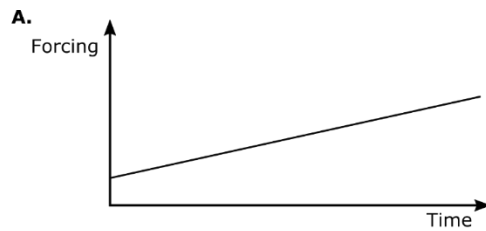
491 The time scales for *process-focused models* are defined by equations (14) and (15) where
492 the latter formulation by Tsujimoto (1990) was derived by considering the Manning-equation,
493 i.e. by considering additional similarity in bed roughness. Time scales for models with
494 suspended load were summarized by e.g. Hughes (1993) and Van Rijn *et al.* (2011), but in
495 almost all cases a morphological time scale of suspended models was derived
496 corresponding to $t_{\eta r} = h_r^{0.5}$ (where the vertical length scale characterizes wave
497 characteristics). These similarity conditions can result in rather impractical scaling ratios,
498 especially when considering both vertical and horizontal directions, and result in a challenge
499 in developing strictly scaled models containing both sediment and water.

500 **3.2. Representing variable forcing and sequences of events**

501 Future climate regimes are anticipated to be characterised by increased variability and higher
502 frequency and magnitude of extreme events such as river flooding (Table 1, Figure 5). Due to the
503 difficulties in scaling unsteady flows and sediment transport in physical models (see section 3.1),
504 there are few physical modelling studies exploring sequences of multiple floods (e.g. Braudrick *et al.*,
505 2009). In terms of improving our understanding of the impact of climate change on fluvial
506 environments, it would be particularly relevant to investigate variations in hydrograph characteristics
507 (i.e. duration, magnitude and frequency) over time scales that are similar to the system recovery time

508 for morphodynamics and vegetation. All systems have a characteristic time scale for recovery
509 following a perturbation (Brunsdon and Thornes 1979). This time scale can range from $>10^3$ years in
510 erosive bedrock settings (e.g. canyons; Baynes *et al.* 2015) to $10^1 - 10^2$ years in alluvial depositional
511 fluvial environments (e.g. sandur plains; Duller *et al.* 2014) due to the relative differences in the
512 mobility of sediments, although larger systems typically take longer to fully recover following a
513 perturbation (Paola 2000). This illustrates that the timing of sequences of flood events relative to the
514 time scale of recovery is as important in driving evolution and change in fluvial environments as the
515 magnitude of individual flood events (Figure 5). With an increased frequency of extreme events, this
516 recovery timescale may be threatened, with subsequent events of possibly greater magnitude
517 occurring before the system has fully recovered from the initial perturbation with potentially unknown
518 consequences. Thus, the accurate representation of non-constant forcing and the relative importance
519 of sequences of events within physical models remains an important goal for the development of the
520 understanding of fluvial system response to future climate scenarios.

521 Traditionally, flood events are represented in physical models at the event scale by triangular
522 hydrographs with possibly an asymmetry between the rising and falling stages (e.g Lee, Liu
523 and Cheng, 2004). The gradual increase and decrease of discharge are reproduced by
524 stepped hydrographs with the number of steps for each hydrograph strongly dependent on
525 the complexity of the flume control equipment (Lee *et al.* 2004; Ahanger *et al.* 2008).
526 Sequences of flood events modelled on a particular system, or the long-term evolution of a
527 system driven by a long-term shift in the magnitude or frequency of forcing are rarely
528 represented in physical models (Figure 5).



529

530 Figure 5. Conceptual diagram indicating different forcing regimes in fluvial and fluvially-affected
 531 systems such as river deltas and estuaries under climate change. (A) A progressive increase in a
 532 constant forcing over a long time scale (e.g. sea level rise, or increase in biostabilisation as a result
 533 of temperature increase). (B) A forcing regime characterised by infrequent and low-magnitude
 534 extreme events, superimposed on the progressive trend shown in (A). (C) A forcing regime
 535 characterised by higher magnitude extreme events, but of the same frequency, compared to (B). (D)
 536 A forcing regime characterised by extreme events of the same magnitude as (B), but occurring more
 537 frequently. (E) A forcing regime characterised by extreme events that are both more frequent and of
 538 a higher magnitude compared to (B). The typical time for the system to recover back to equilibrium
 539 conditions is shown in grey in B-E. Due to frequency and magnitude of the extreme events in (E),
 540 the system has never fully recovered before the subsequent extreme event, placing the system in a
 541 constant state of transience.

542

543 3.3. Representing biology and timescales of biological change

544 Currently, most hydraulic facilities are not well suited to work with living organisms. These facilities
 545 may therefore result in biota being stressed by one or more environmental factors including
 546 inappropriate water chemistry (salinity, pH, dissolved oxygen, inorganic carbon), water temperature,
 547 substrate (physical and chemical properties, soil saturation), lighting (composition, intensity, timing),
 548 and flow characteristics (depth, velocity, drag). The health and behaviour of living plants may also
 549 be affected by biological considerations, including insufficient nourishment (type, quantity, and
 550 28

550 timing), competition for resources amongst individuals and, potentially, the introduction of pathogens.
551 Johnson *et al.*, (2014a) provide a review of these main stressors and their management in flume
552 facilities. Of course, plants are often stressed in their natural environment by competition for
553 resources and by other ecological and biological interactions. Their interactions with their
554 environment are variable and complex, such that there is no ideal stress-free state that must be
555 mimicked. Nevertheless, a basic goal of most experimental work will be to reproduce in the flume
556 behaviours that are typical in nature and, in that case, low levels of stress are desirable, or the
557 development of surrogates that accurately replicate plant/microbial activity and can be time scaled.

558 Most plants are able to tolerate a range of environmental conditions, with fatality beyond limiting
559 thresholds. As conditions become less optimal, but sub-lethal, the plant will adapt, potentially altering
560 the way in which it interacts with the flow. We know very little about these adaptations and what they
561 mean for hydraulic performance, but existing work suggests that the relations are likely to be complex,
562 especially where multiple stressors are present (Puijalon *et al.* 2007).

563 Demonstrating that vegetation is not physiologically or behaviourally stressed during
564 experiments should be a standard element of any physical modelling experiment involving
565 live plants. Without that assurance it is difficult to be confident that measured hydraulic and
566 morphodynamic responses can be properly assigned to treatment effects, not abnormal
567 behaviour caused by the physical modelling environment. While it may be relatively easy to
568 detect serious ill-health or the death of a plant that is part of a flume experiment, earlier
569 stages of decline that affect the plants interaction with the flow, may go undetected,
570 potentially undermining the results obtained.

571 This leads to the identification of two key challenges for investigating plant-flow-sediment
572 interactions: i) developing protocols that can be used to monitor plant health or stress levels
573 during physical modelling experiments, and ii) developing a fuller understanding of how
574 health and stress levels affect key plant structures, physiological responses and behaviours

575 that are relevant to flow and sediment interactions. Meeting these challenges would provide
576 a basis for making objective decisions about how stressed a plant is and whether the level
577 of stress is sufficient to affect its biomechanical behaviour as that affects its interactions with
578 the flow and therefore the integrity of an experiment.

579 From a scaling perspective, of primary interest is the role of the hydraulics as a driving force
580 for the growth and, hence, the geometrical and mechanical properties of plants and biofilms.
581 Hydrological modifications, driven by climate change, especially in terms of flood intensity
582 and frequency, are very likely to also modify plant diversity and distribution (Garssen *et al.*
583 2015). Importantly, the time scales associated with plant and biofilm growth in the field are
584 very large when compared to the time scales of physical modelling experiments in the
585 laboratory. For photosynthetic biofilms in rivers, for example, growth cycles are associated
586 with time scales of around 30 days, which correlates approximately to inter-flood periods in
587 the field (see e.g. Boulêtreau *et al.*, (2010)). Macrophytes or riparian vegetation generally
588 develop and grow over much longer time scales. For biofilms, another issue is the extreme
589 versatility of this biological agent, whose growth and composition adapts very quickly to flow
590 conditions during growth; for example, Graba *et al.*, (2013) demonstrated that in steady-flow
591 growth experiments the biofilms optimized their mechanical properties to fit the imposed
592 steady forcing, and were very easily detached by a slight increase of flow velocity.
593 Incorporating flow unsteadiness associated with typical discharge fluctuations then
594 becomes important for growing representative laboratory biofilms.

595 Plants and biofilms can be simplified and represented by some physical or chemical
596 surrogates. As far as plants are concerned, the use of physical surrogates offers the
597 opportunity to better control the interactions between aquatic vegetation and a changing
598 hydraulic environment, without the issue of phenotypic plasticity typical from biotic systems
599 (Read & Stokes 2006; Nikora 2010). However, the development of surrogates relies on the
600 good understanding of the plant biomechanical properties and requires therefore extensive
30

601 field data collection prior to the main experiments (Nikora 2010). Although recent works are
602 relying more and more on plant surrogates (see Johnson *et al.*, (2014b) for a non-exhaustive
603 list), only a few studies investigated the surrogate design process for complex shaped
604 aquatic plants, such as the work carried out by Paul and Henry (2014), and this process is
605 yet to be developed for freshwater aquatic vegetation.

606

607 **4. Innovative approaches and required future developments to represent** 608 **climate change impacts in physical models.**

609 **4.1. Bridging the timescale gap**

610 The range of physical modelling approaches highlighted in Figure 2 have worked well for both small
611 and large spatial and temporal scales. At the event scale, 1:1 physical models have proven invaluable
612 tools to examine the effects of storm wave on flooding risk and safety (Figure 3). More extreme
613 storm wave and river flood events are projected as a result of climate change (Table 1). The current
614 hydraulic facilities are however expected to incorporate these more extreme events in their
615 experiments seamlessly by adjusting their test scenarios to include the latest climate projections (e.g.,
616 wave height). Other than potentially running into size limitations of the hydraulic facility (i.e. larger
617 events require larger facilities for 1:1 modelling), these more extreme events do not require additional
618 scaling compared to default extreme event tests. This observation indicates that no problems are
619 foreseen in representing more extreme events associated with climate change in hydraulic facilities.

620 Also at larger spatial (landscapes) and temporal ($>10^2$ years) scales, *analogue* models have worked
621 well leading to agenda-setting research and understanding of landscape evolution processes
622 (Hasbargen & Paola 2000; Turowski *et al.*, 2006; Tal & Paola 2007; Bonnet 2009). *Analogue* models
623 can act as a tool for exploration, due to the ability to simplify aspects of a complicated system and
624 explore the behaviour of targeted processes under controlled conditions (Bonnet 2009). The freedom

625 given by foregoing the strict scaling laws can potentially allow innovative experiments to develop an
626 understanding of systems that are manipulated in ways that would not be possible using a strict scaling
627 approach, such as coastal dynamics and response to sea-level rise (Kim *et al.* 2006) or the exploration
628 of different sequences of events on the overall system behaviour (e.g., Ganti *et al.* 2011). It is
629 important to note that *analogue models* are exclusively fit for these ‘thought-provoking’ experiments
630 and hence our primary tool for investigating processes, interactions and feedbacks across longer ($>10^2$
631 years) time scales relevant for climate adaptation purposes (Figure 2).

632 Intermediate time scales ($10^1 - 10^2$ years) have proven difficult to represent in physical models to date,
633 leaving us with a timescale gap in physical modelling capabilities. Yet, in the context of climate
634 change adaptation for planning and policy purposes, the evolution of fluvial systems due to climate
635 change over intermediate time scales is most prevalent and urgent (Figure 2). Depending on the exact
636 timescale or process of interest, *undistorted*, *distorted* and *process-focused* models may provide
637 physical scaling approaches to study the fluvial system at hand. *Undistorted* and *distorted scaled*
638 *models* are best suited to investigate individual and short-lived events due to the minimum
639 compression of spatial and temporal scales (Figure 2). As a consequence, these types of physical
640 models should be used to further our understanding of the projected changes in magnitude and
641 frequency of short-lived events and process interactions, extending the individual event scale covered
642 by 1:1 models. Similarly, *process-focused* and perhaps some distorted and analogue-reach physical
643 models are best placed to condense the timescales represented in analogue models in an effort to study
644 the effects of intermediate timescales of climate change in fluvial systems (Figure 2). At these
645 interdecadal timescales, the effects of variable forcing, sequences of events and biological
646 interactions are dominant (Garssen *et al.* 2014; Garssen *et al.* 2015) but poorly understood drivers of
647 fluvial system behaviour. Bridging the intermediate timescale gap will thus require an effort to
648 implement these drivers in the appropriate physical models for researchers to be able to study the
649 effects of climate change across intermediate timescales. Below, we provide examples of studies on
650 variable forcing, sequences of events, lightweight sediment and biology and we discuss how they can

651 be applied to better represent climate change at intermediate timescales specifically and expand the
652 future physical modelling capability more generally.

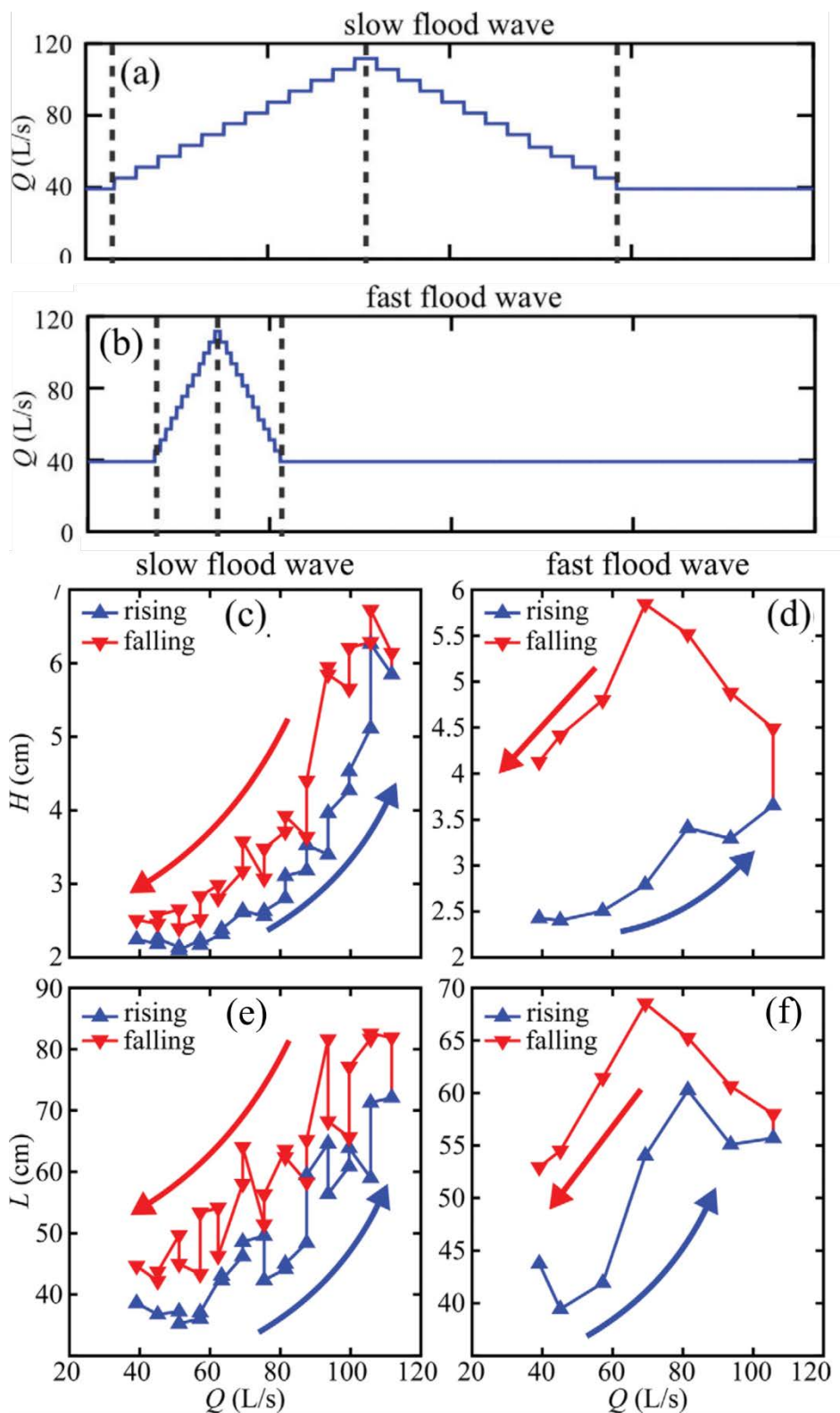
653

654 **4.2. Variable forcing and event sequences**

655 Recently, Martin and Jerolmack (2013) have advanced the the knowledge of bedform
656 dynamics for non-stationary flows, including the difference in the scaling of morphodynamic
657 and hydrodynamic processes (Section 3.1). The processes associated with the growth of
658 bedforms following an abrupt increase in discharge and their decay following an abrupt
659 decrease in discharge are complex and very different (Martin and Jerolmack, 2013). The
660 former relies on gradual collision and merging of small structures towards larger ones, while
661 the latter relies on the formation of secondary small scale structures that cannibalize
662 progressively the large structures formed earlier during the rising stage (Martin and
663 Jerolmack 2013). The timescale of the bedform response under these conditions is
664 proportional to the reconstitution time, defined as $T_r = V/q_s$ where V is the volumetric
665 sediment displacement for the bedform adjustment and q_s is the sediment flux (Martin and
666 Jerolmack, 2013). The reconstitution time is a function of the equilibrium bedform heights,
667 and celerities under the initial and secondary discharge magnitudes, such that taller and
668 longer bedforms take longer to return to equilibrium following an abrupt change in discharge.

669 Additionally, the mechanism and characteristics of the forcing change (i.e., discharge) was
670 found to be important in setting the mechanism of bedform response on the channel bed
671 (Figure 6). Dependent on the rate of a gradual increase and decrease in the discharge
672 (Figure 6a-b), bedforms either respond through a phase of hysteresis or through a linear
673 response of the length and height (Figure 6c-f). Under the ‘fast flood wave’ conditions, the
674 timescale response of the bedform adjustment is shorter than the timescale of flood wave
675 discharge, forcing the hysteresis response. These observations following their experiments

676 under variable forcing allowed Martin and Jerolmack (2013) to propose a simple model
677 framework for the quantitative prediction of bedform adjustment timescale and the
678 occurrence of bedform hysteresis in natural rivers during individual or sequences of events.
679 This innovative example demonstrates the future potential for physical models in advancing
680 the understanding of the processes and response of fluvial systems under variable forcing
681 conditions, aiding the understanding of the possible impacts of climate change. The
682 identification of response timescales of morphodynamic processes to individual events (i.e.,
683 Martin and Jerolmack, 2013) can act as a starting point for evaluating the response to
684 sequences of multiple events of different frequencies and magnitudes (Figure 6).



685

686 Figure 6. Comparison of bedform dynamics under different variable discharge regimes. (a)
 687 Hydrograph simulating a slow flood wave. (b) Hydrograph simulating fast flood wave. (c-d) Evolution
 688 of bedform height during the hydrographs. (e-f) Evolution of bedform length during the hydrographs.
 689 A clear hysteresis is apparent in the evolution of the bedforms during the fast flood wave, due to time
 690 lag of response of the bedforms is greater than the timescales of the flood waves. Adapted from
 691 Martin and Jerolmack (2013).

35

692 Differences in sediment transport rate for different magnitude and duration flood events would not
693 necessarily mean that the order of events would not be important for experiments investigating
694 sequences of events. However, the reorganisation of bed morphology either in terms of bedform size
695 or bed structure through events will impact on the state of the system for the next event, which means
696 that the order of events could be significant and this should be addressed in flume experiments that
697 investigate longer time scales.

698 **4.3. Lightweight sediments**

699 Lightweight materials have been used to study local erosion processes such as scour development
700 downstream of weir structures (e.g., Ettmer, 2006, and references therein), bridge piers and
701 abutments (Fael *et al.*, 2006; Ettmer *et al.*, 2015) and the impact of jets (e.g. Rajaratnam & Mazurek
702 2002). The latter studies, in particular, made use of the fact that erosion processes are accelerated
703 when lightweight sediments are used instead of natural fluvial sediments, i.e. that the equilibrium
704 dimensions of the scour can be reached faster, allowing the time scales of study to be extended
705 (Figure 2). At a larger scale, Willson *et al.*, (2007) reported on a distorted scale model focusing
706 on river and sediment diversions in the lower Mississippi river delta with $L_r = 1:12,000$ and
707 $h_r = 1:500$ and a model sediment with a density $\rho_s = 1050 \text{ kg/m}^3$ covering 77 river miles and
708 an area of about 3526 square miles. In this model, the flow was scaled via the Froude law
709 and the lightweight sediment was scaled based on considerations for the incipient motion of
710 the particles so that incipient motion and resuspension were similar in model and prototype.
711 The resultant sediment time scale was given by the authors with 1:17,857 (one year of
712 prototype time equals roughly 30 minutes of model time). This model was run for different
713 scenarios, including sea-level rise, and used to enhance the general understanding of the
714 impact of planned measures for US State and Federal Agencies (Willson *et al.*, 2007). Such
715 approaches, specifically using lightweight sediment to reduce the time scale of the
716 environmental processes in the physical models can extend the timescale of scaled models

717 (Figure 2) to bridge the gap in modelling capabilities over the timescale relevant for climate
718 change.

719 **4.4. Representing biology**

720 Time scales associated with the growth and behaviour of vegetation are inherently difficult
721 to downscale in physical models using *undistorted* or *distorted* models. Therefore, it is more
722 convenient to use living or artificial surrogates within the analogue modelling approach,
723 where the effects of vegetation in the system are replicated, but not necessarily directly.
724 Plant surrogates also offer new possibilities to test hypotheses in the context of changing
725 fluvial systems. Johnson *et al.*, (2014b) detailed the various benefits and the limitations of
726 using inert physical surrogates, and these points will therefore not be detailed here. Yet,
727 surrogate development is still in its infancy and depends on a detailed knowledge of the
728 morphology and biomechanics of the species of interest, and we present here some of the
729 major issues yet to be tackled, in the context of changing fluvial systems.

730 The morphology and mechanics of aquatic plants can vary based on seasonal patterns. In
731 flume experiments, the potential interaction between the different time scales such as the
732 seasonal growth and the time between active and inactive hydrological regimes needs to be
733 considered. In the case of experiments involving time compression (*analogue* or *process-*
734 *focused models* always active/in flood, see e.g. Paola (2000)) effects due to seasonal
735 changes of plant characteristics may be lost. A good understanding of the plant
736 biomechanical properties requires the use of a solid dataset from real-life conditions (Nikora
737 2010), collected using well identified techniques (Henry 2014; Henry, 2018). Additionally,
738 the required level of complexity of a plant surrogate is still uncertain, as it is critical not to
739 simply redesign the plant structure (Denny 1988). The understanding of the existing
740 structural organisation of a plant is key to the identification of the environmental factor that
741 defined it, and should highlight the features to be reproduced in an experiment, depending

742 on the processes and scales to be investigated. The most important part in a design process,
743 i.e. performance tests, should be conducted systematically to ensure that the dynamic
744 behaviours of the surrogate correspond to the original criteria, i.e. the reproduction of the
745 process observed in nature (flexibility, plant to plant interaction, effect on sediment transport).
746 The application of models without scaling to address questions relating to climate change
747 has some limitations because model time is no faster than prototype time, but for
748 understanding some interactions between organisms and their surroundings, there are no
749 satisfactory scaling relationships (e.g. Wilcock *et al.*, 2008). Kui *et al.*, (2014) present results
750 from the StreamLab experiments that are used to elucidate the ecogeomorphic feedbacks
751 between riparian tree seedlings and flood events. These 1:1 physical models investigate the
752 use of flood releases to control invasive vegetation, however this type of model has the
753 potential to improve our understanding of the response of trees and other organisms to
754 extreme events that could be associated with climate change.

755 In theory, it is possible to scale down plant properties within the distorted scale modelling
756 approach, which may lead to a distortion in time and/or space of the hydraulic model
757 (Johnson *et al.*, 2014b). In practice, no such work has been published to the best of our
758 knowledge, and investigations related to scaled plant properties are just about to start. The
759 interaction of this new distorted 'plant time scale' with the other time scales applying to
760 sediment transport and larger morphological evolutions, is yet to be characterised but offers
761 a potentially important avenue for future work into the holistic evolution of river systems
762 under climate change forcing.

763 For plants, several studies have relied on the use of *alfalfa* because of its size and growth
764 time scale fit with a downscaling approach to physical modelling of sediment and flow
765 dynamics and their interactions with vegetation. This analogue modelling approach leads to
766 floodplains vegetated by a single species that resembles a very fast growing tree (Figure 7).

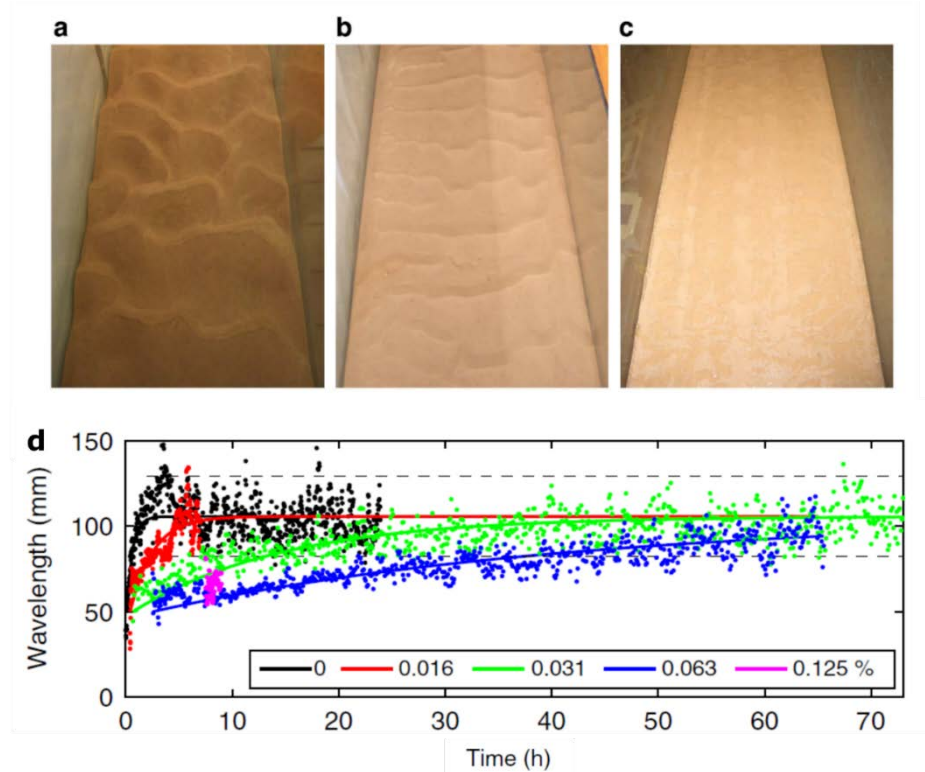
767 Vegetation is able to stabilise river banks, focus and organise the flow and hence convert the planform
768 morphology from braided to single-thread (Gran and Paola 2001; Tal *et al.* 2004; Tal and Paola 2007;
769 Tal and Paola 2010; Braudrick *et al.* 2009; Van de Lageweg *et al.* 2010; Van Dijk *et al.* 2013; Bertoldi
770 *et al.* 2014). It should be noted, however, that vegetation alone does not lead to fully meandering
771 channels (Desloges and Church, 1989) and fine grained material is also required (Santos *et al.*,
772 2017a,b). Morphological trends associated with the colonisation of a floodplain by riparian vegetation
773 are an increased sinuosity, lower lateral migration rates, a reduced number of channels, deepening of
774 the channels, and a reduction in the wetted area, and potentially can provide insights into the large-
775 scale evolution of river systems under climate-induced variability into vegetation patterns.



776
777 Figure 7. Example of a physical model in which the original fluvial braided plain has been colonised
778 by small-scale alfalfa vegetation. Flow is from right to left and the panel is 6 m long and 2 m wide.

779 In addition to plant surrogates, it may be possible to use chemical surrogates to simulate
780 aspects of biofilm mediated stabilization processes. Xanthan gum (a rheology modifier often
781 used in the food industry) is one example of such a surrogate and has been employed in a
782 number of studies to mimic natural biofilm behaviour (Black *et al.* 2001; Tolhurst *et al.* 2002).
783 Even though it has been demonstrated that Xanthan gum is not a perfect analogue of natural
784 biofilms (Perkins *et al.* 2004), primarily because natural biofilms are more complex, it is
785 seemingly useful in studies on sediment erosion, with increasing quantities of Xanthan gum
786 having a clear effect on the morphology of bedforms (Malarkey *et al.*, 2015; Schindler *et al.*,
787 2015; Parsons *et al.*, 2016; Figure 8). A recent experimental investigation compared the
788 stabilisation effects for sand of Xanthan Gum to three other chemical surrogates; Alginic
789 Acid, Carrageenan and Agar (van de Lageweg *et al.*, 2017). Alginic Acid and Agar had a
39

790 limited effect, as the erosion threshold for the sediment did not increase while the erosion
791 threshold increased linearly for increased concentrations of Xantham Gum and
792 Carrageenan (van de Lageweg et al., 2017), potentially providing a method of speeding up
793 time scales of physical modelling experiments investigating biostabilization effects.



794

795 Figure 8. Effect of extracted extracellular polymeric substances (EPS) content on bedform
796 morphology. (a) 0% EPS, (b) 0.125% EPS content, (c) 1% EPS content. (d) Ripple height
797 development for different EPS contents (black: 0%, red: 0.016%, green: 0.031%, blue: 0.063%, pink:
798 0.125%). Adapted from Malarkey *et al.*, (2015).

799

800 4.5. Infrastructural developments

801 A potential barrier preventing the implementation of the innovative approaches discussed above are
802 the physical limitations of the infrastructure associated with the available physical modelling facilities.

803 An obvious example, given the potential stresses placed on growing plants and vegetation in the
804 unnatural conditions of many physical modelling laboratories, is improved facilities designed for
805 optimal biological growth. Potential developments include climate and light-controlled conditions,

806 nutrient delivery, and stress monitoring protocols during the set up and duration of experiments
807 (Johnson *et al.*, 2014a). An additional infrastructural development that is required relates to the
808 measurement and monitoring techniques employed during physical modelling experiments.
809 Especially as the understanding of the impact of climate change and variable forcing in fluvial
810 systems requires a quantification of both short-term and longer-term dynamics (e.g. the impact of
811 single storm events on top of the longer term impact of gradual sea-level rise). Monitoring and
812 measuring remains a challenge for studies that aim to quantify and disentangle the impact from
813 individual short-lived events to longer-term trends due to the lack of high resolution monitoring and
814 quantification techniques that can operate over multiple time scales (Kim *et al.* 2006). It is
815 recommended that future studies investigate deltaic and estuarine environments with combined
816 fluvial and tidal currents, and the Metronome tidal facility at the University of Utrecht is an innovative
817 facility that has been developed in recent years (Kleinhans *et al.* 2017). These experiments could
818 provide the ability to observe, monitor and characterise the driving processes that lead to the transition
819 between different equilibrium conditions, and the balance of different aspects of the fluvial
820 landscapes and ecosystems in tidally-dominated environments. This would also improve the
821 parameterisation of such processes in numerical models and associated predictions of how fluvial
822 systems may respond to variations in climatic forcing.

823 **4.6. Linkages with numerical simulations**

824 It is anticipated that combining physical modelling and numerical modelling has the potential to be a
825 robust way forward to address the current gap in the capability to model climate change adaptation.
826 For example, physical modelling can be used to perform focussed sensitivity analyses on the impact
827 of individual parameters in controlled environments, aiding the parameterisation of numerical models
828 that simulate processes such as flow-vegetation interactions (Marjoribanks *et al.* 2015). Numerical
829 models parameterised from empirical data have explored scenarios and provided projections for the
830 evolution of fluvial landscapes (Coulthard, Hicks, and Van De Wiel 2007; Nicholas and Quine 2007;
831 Attal *et al.* 2008; Nicholas 2013; Edmonds and Slingerland 2010; Schuurman *et al.* 2013; Liang *et al.*
41

832 2016) and sediment-vegetation interactions in these systems (van Oorschot *et al.* 2016). It may be
833 noted that the development of the use of inert plant surrogates may also help and be done
834 in parallel to numerical modelling studies replicating fluid flow around vegetation
835 (Marjoribanks *et al.* 2014; Marjoribanks, Hardy, and Parsons 2015), whose effects can be
836 included into larger numerical simulation addressing fluvial adaptation at a larger space and
837 time scale.

838 Numerical models can be used to explore which combinations of variables are most worth studying
839 in physical experiments and can aid with the planning of such experiments. Once accurately
840 parameterised and calibrated in physical models, process-based numerical models could be upscaled
841 to cover larger spatial scales and longer time periods that are appropriate for climate change
842 adaptation (i.e. intermediate scales). Also, numerical model simulations can be useful predictive
843 tools because they can cover multiple spatial and temporal scales and they can easily be forced with
844 a multitude of climate change scenarios that would be impractical using physical models. However,
845 these numerical simulations often contain associated uncertainty due to the inability to determine
846 whether the observed behaviour is a result of true landscape dynamics or merely an artefact of the
847 model set up. Physical models could potentially improve this confidence by replicating some of the
848 same scenarios and comparing the behaviour and interactions between processes in both the
849 numerical and physical simulations.

850

851 **5. Conclusions**

852 Physical modelling has contributed significantly to our understanding of fluvial systems. This is
853 expected to continue into the future as different physical modelling approaches are well suited to
854 investigate the response and potential adaptation to climatically driven changes in forcing over
855 various timescales. Based on a review of the state-of-the-art in physical modelling of fluvial systems,
856 this study highlights that: (i) physical modelling offers a prime opportunity for furthering the current

857 understanding of variability of forcing in fluvial systems. (ii) For the study of fluvial systems for
858 climate change adaptation for policy purposes, the modelled time scales using *1:1, undistorted* or
859 *distorted scale models* need to be extended and the modelled time scales using *process-focused* or
860 *analogue* need to be reduced in order to bridge a gap over the relevant decadal timescales. (iii)
861 Representing the response of plants and organisms to changing conditions and the resulting feedback
862 on physical processes requires more attention and better techniques than presently available, using
863 both *distorted scale* and *analogue* surrogate modelling approaches. (iv) Coupling of physical
864 modelling output with numerical model parameterisation and development is crucial for producing
865 accurate predictions of how fluvial systems will respond in the future to a range of possible forcing
866 scenarios over multiple time scales.

867 Within the context of climatic change in fluvial environments, future focus and investment is
868 recommended towards the physical modelling of the detailed interactions between riverine biology,
869 hydrology and morphology, non-constant forcing and an understanding of the impacts of single
870 events, multi-decadal oscillations and longer term trends. This will enable the development of
871 appropriate and effective mitigation strategies for fluvial ecosystems and environments under threat
872 from climate change, that are grounded in robust physical experimentation.

873

874 **6. Acknowledgments**

875 We thank Jeff Peakall and one anonymous reviewer for the constructive comments on an
876 earlier version of this manuscript. The work described in this publication was supported by
877 the European Community's Horizon 2020 Programme through the grant to the budget of the
878 Integrated Infrastructure Initiative HYDRALAB+, Contract no. 654110.

879

880 7. References

- 881 Ahanger, M.A., G.L. Asawa, and M.A. Lone. 2008. "Experimental Study of Sediment Transport
882 Hysteresis." *Journal of Hydraulic Research* 46 (5). Taylor & Francis Group: 628–35.
883 doi:10.3826/jhr.2008.3185.
- 884 Ashmore, P E. 1988. "Bed Load Transport in Braided Gravel-Bed Stream Models." *Earth Surface
885 Processes and Landforms* 13: 677–95.
- 886 Attal, M., G.E. Tucker, A.C. Whittaker, P.A. Cowie, and G.P. Roberts. 2008. "Modeling Fluvial
887 Incision and Transient Landscape Evolution: Influence of Dynamic Channel Adjustment."
888 *Journal of Geophysical Research* 113 (F3): F03013. doi:10.1029/2007JF000893.
- 889 Barenblatt, G.I. 2003. *Scaling*. Vol. 34. Cambridge University Press.
- 890 Baynes, E.R.C., M. Attal, S. Niedermann, L.A. Kirstein, A.J. Dugmore, and M. Naylor. 2015.
891 "Erosion during Extreme Flood Events Dominates Holocene Canyon Evolution in Northeast
892 Iceland." *Proceedings of the National Academy of Sciences of the United States of America*
893 112 (8). National Academy of Sciences: 2355–60. doi:10.1073/pnas.1415443112.
- 894 Bennett, S.J., P. Ashmore, and C. McKenna Neuman. 2015. "Transformative Geomorphic Research
895 Using Laboratory Experimentation." *Geomorphology* 244 (September). Elsevier: 1–8.
896 doi:10.1016/J.GEOMORPH.2014.11.002.
- 897 Bertoldi, W., M. Welber, L. hav, S. Zanella, and F. Comiti. 2014. "A Flume Experiment on Wood
898 Storage and Remobilization in Braided River Systems." *Earth Surface Processes and
899 Landforms* 39 (6): 804–13. doi:10.1002/esp.3537.
- 900 Best, J L. 1998. "The Influence of Particle Rotation on Wake Stability at Particle Reynolds
901 Numbers, $Re P < 300$ —Implications for Turbulence Modulation in Two-Phase Flows."
902 *International Journal of Multiphase Flow* 24 (5). Elsevier: 693–720.

- 903 Black, K.S., H. Sun, G. Craig, D.M. Paterson, J. Watson, and T.J. Tolhurst. 2001. "Incipient
904 Erosion of Biostabilized Sediments Examined Using Particle-Field Optical Holography."
905 American Chemical Society. doi:10.1021/ES0014739.
- 906 Bonnet, S. 2009. "Shrinking and Splitting of Drainage Basins in Orogenic Landscapes from the
907 Migration of the Main Drainage Divide." *Nature Geoscience* 2 (11). Nature Publishing Group:
908 766–71. doi:10.1038/ngeo666.
- 909 Boulêtreau, S., M. Sellali, A. Elozegi, Y. Nicaise, Y. Bercovitz, F. Moulin, O. Eiff, S. Sauvage, J.
910 Sánchez-Pérez, and F. Garabétian. 2010. "Temporal Dynamics of River Biofilm in Constant
911 Flows: A Case Study in a Riverside Laboratory Flume." *International Review of Hydrobiology*
912 95 (2). Wiley Online Library: 156–70.
- 913 Braudrick, C A, W E Dietrich, G T Leverich, and L S Sklar. 2009. "Experimental Evidence for the
914 Conditions Necessary to Sustain Meandering in Coarse-Bedded Rivers." *PNAS* 106 (16): 936--
915 941.
- 916 Brunsden, D., and J.B. Thornes. 1979. "Landscape Sensitivity and Change." *Transactions of the*
917 *Institute of British Geographers* 4 (4). The Royal Geographical Society (with the Institute of
918 British Geographers): 463–84. doi:10.2307/622210.
- 919 Bruun, P. 1966. "Model Geology: Prototype and Laboratory Streams." *GSA Bulletin* 77 (9).
920 GeoScienceWorld: 959–74. doi:10.1130/0016-7606(1966)77[959:mgpals]2.0.co;2.
- 921 Buckingham, E. 1914. "On Physically Similar Systems; Illustrations of the Use of Dimensional
922 Equations." *Physical Review* 4 (4): 345–76. doi:10.1103/PhysRev.4.345.
- 923 Claudin, P., and B. Andreotti. 2006. "A Scaling Law for Aeolian Dunes on Mars, Venus, Earth, and
924 for Subaqueous Ripples." *Earth and Planetary Science Letters* 252 (1–2). Elsevier: 30–44.
925 doi:10.1016/J.EPSL.2006.09.004.
- 926 Coleman, S. E., and V. I. Nikora. 2009. "Exner Equation: A Continuum Approximation of a

- 927 Discrete Granular System.” *Water Resources Research* 45 (9). doi:10.1029/2008WR007604.
- 928 Coulthard, T.J., D.M. Hicks, and M.J. Van De Wiel. 2007. “Cellular Modelling of River
929 Catchments and Reaches: Advantages, Limitations and Prospects.” *Geomorphology* 90 (3).
930 Elsevier: 192–207.
- 931 Davinroy, R.D., R.A. Gaines, D.C. Gordon, L.L. Hopkins, D.P. Barretta, J.L. Brown, M.T.
932 Rodgers, E.A.N. Cox, T. Lauth, and B. Krischel. 2012. “Hydraulic Sediment Response
933 Modeling, Replication Accuracy to the River.” In *Technical Paper M53*. U.S. Army Corps of
934 Engineers, St. Louis and Memphis Districts.
- 935 De Vries, M. 1993. “Use of Models for River Problems.” *Studies and Reports on Hydrology*, 51.
936 Unesco.
- 937 de Vries, M, G J Klaassen, and N Struiksmā. 1990. “On the Use of Movable Bed Models for River
938 Problems: A State-of-the-Art.” *International Journal of Sediment Research* 5 (1): 35–47.
- 939 de Wit, M., and J. Stankiewicz. 2006. “Changes in Surface Water Supply Across Africa with
940 Predicted Climate Change.” *Science* 311 (5769): 1917–21. doi:10.1126/science.1119929.
- 941 Denny, M.W. 1988. *Biology and the Mechanics of the Wave-Swept Environment*. Princeton
942 University Press.
- 943 Desloges, J.R. & Church, M. 1989. Wandering gravel-bed rivers. *Canadian Geographer*,
944 33, 360–364.
- 945 Dietrich, W.E., M.C. Palucis, R.M.E. Williams, K.W. Lewis, F. Rivera-Hernandez and D.Y.
946 Sumner. 2017. Fluvial Gravels on Mars. In: *Gravel Bed Rivers: Processes and Disasters*. D.
947 Tsutsumi & B. Laronne (eds.), Wiley, DOI: 10.1002/9781118971437.ch28
- 948 Duller, R.A., N.H. Warner, C. McGonigle, S. De Angelis, A.J. Russell, and N.P. Mountney. 2014.
949 “Landscape Reaction, Response, and Recovery Following the Catastrophic 1918 Katla
950 Jökulhlaup, Southern Iceland.” *Geophysical Research Letters* 41 (12): 4214–21.

- 951 doi:10.1002/2014GL060090.
- 952 Edmonds, D.A., and R.L. Slingerland. 2010. “Significant Effect of Sediment Cohesion on Delta
953 Morphology.” *Nature Geoscience* 3 (2): 105–9. doi:10.1038/ngeo730.
- 954 EEA. 2017. *Climate Change, Impacts and Vulnerability in Europe 2017: An Indicator-Based
955 Report. EEA Report.* doi:10.2800/66071.
- 956 Einstein, H.A., and N. Chien. 1956. “Similarity of Distorted River Models with Movable Beds.”
957 *Transactions of the American Society of Civil Engineers* 121 (1). ASCE: 440–57.
- 958 Ettema, R., and M. Muste. 2004. “Scale Effects in Flume Experiments on Flow around a Spur Dike
959 in Flatbed Channel.” *Journal of Hydraulic Engineering-ASCE* 130 (7): 635–46.
960 doi:10.1061/(ASCE)0733-9429(2004)130:7(635).
- 961 Ettmer, B.. 2006. “Untersuchungen Zu Kolkvorgängen hinter dem unterströmten
962 Schütz.“ *Mitteilungen des Leichtweiß-Instituts für Wasserbau. Technische Universität
963 Braunschweig, Hef 155, pp. 1-156.*
- 964 Ettmer, B., F. Orth, and O. Link. 2015. “Live-Bed Scour at Bridge Piers in a Lightweight
965 Polystyrene Bed.” *Journal of Hydraulic Engineering* 141 (1999): 1–10.
966 doi:10.1061/(ASCE)HY.1943-7900.0001025.
- 967 Fael, C.M.S., G. Simarro-Grande, J. Martín-Vide, and A.H. Cardoso. 2006. “Local Scour at
968 Vertical-Wall Abutments under Clear-Water Flow Conditions.” *Water Resources Research* 42
969 (10). doi:10.1029/2005WR004443.
- 970 Friedkin, J F. 1945. *A Laboratory Study of the Meandering of Alluvial Rivers.* U.S. Waterways
971 Experiment Station, Vicksburg, Mississippi, USA: U.S. Army Corps of Engineers.
- 972 Frostick, L.E., S.J. McLelland, and T.G. Mercer. 2011. *Users Guide to Physical Modelling and
973 Experimentation: Experience of the HYDRALAB Network.* CRC Press.

- 974 Frostick, L.E., R.E. Thomas, M.F. Johnson, S.P. Rice, and S.J. McLelland. 2014. *Users Guide to*
975 *Ecohydraulic Modelling and Experimentation: Experience of the Ecohydraulic Research Team*
976 *(PISCES) of the HYDRALAB Network*. CRC Press.
- 977 Gaines, R A, and R H Smith. 2002. “Micro-Scale Loose-Bed Physical Models.” In *Hydraulic*
978 *Measurements and Experimental Methods 2002*, 1–12.
- 979 Ganti, V., A. J. Chadwick, H. J. Hassenruck-Gudipati, B. M. Fuller, and M. P. Lamb. 2016.
980 “Experimental River Delta Size Set by Multiple Floods and Backwater Hydrodynamics.”
981 *Science Advances* 2 (5). American Association for the Advancement of Science: e1501768–
982 e1501768. doi:10.1126/sciadv.1501768.
- 983 Ganti, V., K.M. Straub, E Foufoula-Georgiou, and C Paola. 2011. “Space-Time Dynamics of
984 Depositional Systems: Experimental Evidence and Theoretical Modeling of Heavy-Tailed
985 Statistics.” *Journal of Geophysical Research* 116: doi:10.1029/2010JF001893.
- 986 Garssen, A.G., A. Baattrup-Pedersen, L.A.C.J Voeselek, J.T.A Verhoeven, and M.B. Soons. 2015.
987 “Riparian Plant Community Responses to Increased Flooding: A Meta-Analysis.” *Global*
988 *Change Biology*. doi:10.1111/gcb.12921.
- 989 Garssen, A.G., J.T.A. Verhoeven, and M.B. Soons. 2014. “Effects of Climate-Induced Increases in
990 Summer Drought on Riparian Plant Species: A Meta-Analysis.” *Freshwater Biology* 59 (5):
991 1052–63. doi:10.1111/fwb.12328.
- 992 Gorrick, S., and J.F. Rodríguez. 2014. “Scaling of Sediment Dynamics in a Laboratory Model of a
993 Sand-Bed Stream.” *Journal of Hydro-Environment Research* 8 (2). Elsevier: 77–87.
994 doi:10.1016/J.JHER.2013.12.001.
- 995 Graba, M., S. Sauvage, F.Y. Moulin, G. Urrea, S. Sabater, and J.M. Sanchez-Pérez. 2013.
996 “Interaction between Local Hydrodynamics and Algal Community in Epilithic Biofilm.”
997 *Water Research* 47 (7): 2153–63. doi:10.1016/j.watres.2013.01.011.

- 998 Gran, K, and C Paola. 2001. "Riparian Vegetation Controls on Braided Stream Dynamics." *Water*
999 *Resources Research* 37 (12): 3275–83.
- 1000 Grass, A.J. 1971. "Structural Features of Turbulent Flow over Smooth and Rough Boundaries."
1001 *Journal of Fluid Mechanics* 50 (2). Cambridge University Press: 233–55.
1002 doi:10.1017/S0022112071002556.
- 1003 Grasso, F., H. Michallet, E. Barthélemy, and R. Certain. 2009. "Physical Modeling of Intermediate
1004 Cross-Shore Beach Morphology: Transients and Equilibrium States." *Journal of Geophysical*
1005 *Research* 114 (C9): C09001. doi:10.1029/2009JC005308.
- 1006 Green, D.L. 2014. "Modelling Geomorphic Systems: Scaled Physical Models." In
1007 *Geomorphological Techniques. British Society for Geomorphology*, edited by J.M. S. J. Cook,
1008 Clarke, L.E. & Nield. London, UK, UK: British Society for Geomorphology. ISSN: 2047-
1009 0371.
- 1010 Guerit, L., F. Métivier, O. Devauchelle, E. Lajeunesse, and L. Barrier. 2014. "Laboratory Alluvial
1011 Fans in One Dimension." *Physical Review E* 90 (2). American Physical Society: 22203.
1012 doi:10.1103/PhysRevE.90.022203.
- 1013 Harley, C.D.G., A.R. Hughes, K.M. Hultgren, B.G. Miner, C.J.B. Sorte, C.S. Thornber, L.F.
1014 Rodriguez, L. Tomanek, and S.L. Williams. 2006. "The Impacts of Climate Change in Coastal
1015 Marine Systems." *Ecology Letters* 9 (2): 228–41. doi:10.1111/j.1461-0248.2005.00871.x.
- 1016 Hasbargen, L. E., and C. Paola. 2000. "Landscape Instability in an Experimental Drainage Basin."
1017 *Geology* 28 (12): 1067–70. doi:10.1130/0091-7613(2000)28<1067:LIAED>2.0.CO.
- 1018 Heller, V. 2011. "Scale Effects in Physical Hydraulic Engineering Models." *Journal of Hydraulic*
1019 *Research* 49 (3). Taylor & Francis: 293–306. doi:10.1080/00221686.2011.578914.
- 1020 Henry, P.T. 2014. "Bending Properties of a Macroalga: Adaptation of Peirce's Cantilever Test for
1021 in Situ Measurements of *Laminaria Digitata* (Laminariaceae)." *American Journal of Botany*

- 1022 101 (6). Botanical Society of America: 1050–55. doi:10.3732/ajb.1400163.
- 1023 Henry, P.T., 2018. "Variability and similarities in the structural properties of two related *Laminaria*
1024 kelp species". *Estuarine, Coastal and Shelf Science* 200, 395-405.
- 1025 Hentschel, B. 2007. "Hydraulische Flussmodelle Mit Beweglicher Sohle." *Wasserbauliches*
1026 *Versuchswesen* 90: 25–46.
- 1027 Hofland, B, J A Battjes, and R Booij. 2005. "Measurement of Fluctuating Pressures on Coarse Bed
1028 Material." *J. of Hydraulic Engineering* 131 (9): 770–81.
- 1029 Hooke, R. 1968. "Model Geology: Prototype and Laboratory Streams: Discussion." *GSA Bulletin*
1030 79 (3). GeoScienceWorld: 391–94. doi:10.1130/0016-7606(1968)79[391:mgpals]2.0.co;2.
- 1031 Hoyal, D.C.J.D., and B.A. Sheets. 2009. "Morphodynamic Evolution of Experimental Cohesive
1032 Deltas." *Journal of Geophysical Research* 114: F02009. doi:10.1029/2007JF000882.
- 1033 Hudson, R. Y. 1979. "Coastal Hydraulic Models. I. Introduction." *Coastal Hydraulic Models,*
1034 *U.S.Army Coastal Engineering Research Center, Special Report 5, 19–23.*
- 1035 Hughes, S. 1993. *Physical Models and Laboratory Techniques in Coastal Engineering. Advanced*
1036 *Series on Ocean Engineering. Vol. 7.* Singapore: World Scientific Publishing.
- 1037 IPCC. 2014. "Climate Change 2014: Synthesis Report. Contribution of Working Groups I, II and III
1038 to the Fifth Assessment Report of the Intergovernmental Panel on Climate Change."
1039 doi:10.1017/CBO9781107415324.004.
- 1040 Jerolmack, D.J. 2009. "Conceptual Framework for Assessing the Response of Delta Channel
1041 Networks to Holocene Sea Level Rise." *Quaternary Science Reviews* 28 (17): 1786–1800.
1042 doi:10.1016/j.quascirev.2009.02.015.
- 1043 Johnson, M.F., S.P. Rice, W.E. Penning, and J.T. Dijkstra. 2014a. "Maintaining the Health and
1044 Behavioural Integrity of Plants and Animals in Experimental Facilities." In *Users Guide to*

- 1045 *Ecohydraulic Modelling and Experimentation: Experience of the Ecohydraulic Research Team*
1046 *(PISCES) of the HYDRALAB Network*, edited by L.E. Frostick, R.E. Thomas, M.F. Johnson,
1047 S.P. Rice, and S.J. McLelland. CRC Press/Balkema, Leiden, The Netherlands.
- 1048 Johnson, M.F., R.E. Thomas, J.T. Dijkstra, M. Paul, W.E. Penning, and S.P. Rice. 2014b. “Using
1049 Surrogates, Including Scaling Issues, in Laboratory Flumes and Basins.” In *Users Guide to*
1050 *Ecohydraulic Modelling and Experimentation: Experience of the Ecohydraulic Research Team*
1051 *(PISCES) of the HYDRALAB Network*, edited by L.E. Frostick, R.E. Thomas, M.F. Johnson,
1052 S.P. Rice, and S.J. McLelland. CRC Press/Balkema, Leiden, The Netherlands.
- 1053 Kamphuis, J.W. 1975. “Friction Factor under Oscillatory Waves.” *Journal of the Waterways,*
1054 *Harbors and Coastal Engineering Division* 101 (2). ASCE: 135–44.
- 1055 Kim, W., C. Paola, J.B. Swenson, and V.R. Voller. 2006. “Shoreline Response to Autogenic
1056 Processes of Sediment Storage and Release in the Fluvial System.” *Journal of Geophysical*
1057 *Research* 111: F04013, doi:10.1029/2006JF000470.
- 1058 Kleinhans, M.G., C. Braudrick, W.M. Van Dijk, W.I. Van de Lageweg, R. Teske, and M. Van
1059 Oorschot. 2015. “Swiftness of Biomorphodynamics in Lilliput- to Giant-Sized Rivers and
1060 Deltas.” *Geomorphology* 244: 56–73. doi:10.1016/j.geomorph.2015.04.022.
- 1061 Kleinhans, M.G., M. van der Vegt, J. Leuven, L. Braat, H. Markies, A. Simmelink, C. Roosendaal,
1062 A. van Eijk, P. Vrijbergen, and M. van Maarseveen. 2017. “Turning the Tide: Comparison of
1063 Tidal Flow by Periodic Sea Level Fluctuation and by Periodic Bed Tilting in Scaled Landscape
1064 Experiments of Estuaries.” *Earth Surface Dynamics* 5 (4). Copernicus GmbH: 731–56.
- 1065 Kleinhans, M.G., W.M. van Dijk, W.I. van de Lageweg, D.C.J.D. Hoyal, H. Markies, M. van
1066 Maarseveen, C. Roosendaal, et al. 2014. “Quantifiable Effectiveness of Experimental Scaling
1067 of River- and Delta Morphodynamics and Stratigraphy.” *Earth Science Reviews* 133: 43–61,
1068 doi:10.1016/j.earscirev.2014.03.001.

- 1069 Kobus, H. 1978. "Wasserbauliches Versuchswesen." *Schriftenreihe des Deutschen Verbandes für*
1070 *Wasserwirtschaft und Kulturbau*, Mitteilungsheft 4, Parey Verlag, Hamburg.
- 1071 Kraal, E R, M van Dijk, G Postma, and M G Kleinhans. 2008. "Martian Stepped-Delta Formation
1072 by Rapid Water Release." *Nature* 451: 973--976. doi:10.1038/nature06615.
- 1073 Kui, L., J.C. Stella, A. Lightbody, and A.C. Wilcox. 2014. "Ecogeomorphic Feedbacks and Flood
1074 Loss of Riparian Tree Seedlings in Meandering Channel Experiments." *Water Resources*
1075 *Research* 50 (12): 9366–84. doi:10.1002/2014WR015719.
- 1076 Lee, K.T., Y. Liu, and K. Cheng. 2004. "Experimental Investigation of Bedload Transport
1077 Processes under Unsteady Flow Conditions." *Hydrological Processes* 18 (13). John Wiley &
1078 Sons, Ltd.: 2439–54. doi:10.1002/hyp.1473.
- 1079 Liang, M., C. Van Dyk, and P. Passalacqua. 2016. "Quantifying the Patterns and Dynamics of River
1080 Deltas under Conditions of Steady Forcing and Relative Sea Level Rise." *Journal of*
1081 *Geophysical Research: Earth Surface* 121 (2): 465–96. doi:10.1002/2015JF003653.
- 1082 Lick, W, and J Gailani. 2004. "Initiation of Movement of Quartz Particles." *Journnal of Hydraulic*
1083 *Engineering* 130: 755–61.
- 1084 Low, H S. 1989. "Effect of Sediment Density on Bed-Load Transport." *J. of Hydraulic Engineering*
1085 115 (1): 124–38.
- 1086 Lowe, J.A., and J.M. Gregory. 2005. "The Effects of Climate Change on Storm Surges around the
1087 United Kingdom." *Philosophical Transactions of the Royal Society A: Mathematical, Physical*
1088 *and Engineering Sciences* 363 (1831): 1313–28. doi:10.1098/rsta.2005.1570.
- 1089 Lu, J., X. Liao, and G. Zhao. 2013. "Experimental Study on Effects of Geometric Distortion upon
1090 Suspended Sediments in Bending Channels." *Sedimentary Geology* 294 (August). Elsevier:
1091 27–36. doi:10.1016/J.SEDGEO.2013.05.006.
- 1092 Maike, P., and P.T. Henry. 2014. "Evaluation of the Use of Surrogate *Laminaria Digitata* in Eco-
52

- 1093 Hydraulic Laboratory Experiments.” *Journal of Hydrodynamics, Ser. B* 26 (3). Elsevier: 374–
1094 83.
- 1095 Malarkey, J., J.H. Baas, J.A. Hope, R.J. Aspden, D.R. Parsons, J. Peakall, D.M. Paterson, et al.
1096 2015. “The Pervasive Role of Biological Cohesion in Bedform Development.” *Nature*
1097 *Communications* 6 (February). Nature Publishing Group: 6257. doi:10.1038/ncomms7257.
- 1098 Marjoribanks, T.I., R.J. Hardy, S.N. Lane, and D.R. Parsons. 2014. “High-Resolution Numerical
1099 Modelling of Flow - Vegetation Interactions.” *Journal of Hydraulic Research* 52 (6): 775–93.
1100 doi:10.1080/00221686.2014.948502.
- 1101 Marjoribanks, T.I., R.J. Hardy, and D.R. Parsons. 2015. “On Validating Predictions of Plant Motion
1102 in Coupled Biomechanical-Flow Models.” *Journal of Hydraulic Research* 53 (6). Taylor &
1103 Francis: 808–13. doi:10.1080/00221686.2015.1110627.
- 1104 Martin, J., B. Sheets, C. Paola, D. Hoyal. 2009. "Influence of steady base-level rise on channel
1105 mobility, shoreline migration, and scaling properties of a cohesive experimental delta".
1106 *Journal of Geophysical Research* 114, F03017.
- 1107 Martin, R.L., and D.J. Jerolmack. 2013. “Origin of Hysteresis in Bedform Response to Unsteady
1108 Flows.” *Water Resources Research* 49 (3): 1314–1333. doi:10.1002/wrcr.20093.
- 1109 Nicholas, A. 2013. “Morphodynamic Diversity of the World’s Largest Rivers.” *Geology* 41 (4).
1110 GeoScienceWorld: 475–78. doi:10.1130/G34016.1.
- 1111 Nicholas, A., and T.A. Quine. 2007. “Crossing the Divide: Representation of Channels and
1112 Processes in Reduced-Complexity River Models at Reach and Landscape Scales.”
1113 *Geomorphology* 90 (3). Elsevier: 318–39.
- 1114 Nijssen, B., G.M. O’Donnell, A.F. Hamlet, and D.P. Lettenmaier. 2001. “Hydrologic Sensitivity of
1115 Global Rivers to Climate Change.” *Climatic Change* 50 (1/2). Kluwer Academic Publishers:
1116 143–75. doi:10.1023/A:1010616428763.

- 1117 Nikora, V.I. 2010. "Hydrodynamics of Aquatic Ecosystems : An Interface between Ecology,
1118 Biomechanics and Environmental Fluid Mechanics." *River Research and Applications* 26 (4).
1119 John Wiley & Sons, Ltd.: 367–84. doi:10.1002/rra.1291.
- 1120 Nikora, V.I., D.G. Goring, I. McEwan, and G.A. Griffiths. 2001. "Spatially Averaged Open-
1121 Channel Flow over Rough Bed." *J. of Hydraulic Engineering* 127 (2): 123–33.
- 1122 Novak, P., V. Guinot, J. Alan, and D.E. Reeve. 2010. *Hydraulic Modelling: An Introduction,*
1123 *Principles, Methods and Applications.* Spon Press New York
- 1124 Paola, C. 2000. "Quantitative Models of Sedimentary Basin Filling." *Sedimentology* 47: 121–78.
- 1125 Paola, C., and V. R. Voller. 2005. "A Generalized Exner Equation for Sediment Mass Balance."
1126 *Journal of Geophysical Research: Earth Surface* 110 (F4). doi:10.1029/2004JF000274.
- 1127 Paola, C, K M Straub, D C Mohrig, and L Reinhardt. 2009. "The ``unreasonable Effectiveness`` of
1128 Stratigraphic and Geomorphic Experiments." *Earth-Science Reviews* 97: 1–43.
- 1129 Parsons, D.R., R.J. Schindler, J.A. Hope, J. Malarkey, J.H. Baas, J. Peakall, A.J. Manning, et al.
1130 2016. "The Role of Biophysical Cohesion on Subaqueous Bed Form Size." *Geophysical*
1131 *Research Letters* 43 (4): 1566–73. doi:10.1002/2016GL067667.
- 1132 Peakall, J., and J. Warburton. 1996. "Surface Tension in Small Hydraulic River Models - The
1133 Significance of the Weber Number." *Journal of Hydrology New Zealand* 35 (2): 199–212.
- 1134 Peakall, J., P.J. Ashworth, and J. L. Best. 1996. "Physical Modelling in Fluvial Geomorphology:
1135 Principles, Applications and Unresolved Issues." In *The Scientific Nature of Geomorphology,*
1136 edited by B L Rhoads and C E Thorn, 221–53. Chichester, UK: Wiley.
- 1137 Peakall, J, P J Ashworth, and J L Best. 2007. "Meander-Bend Evolution, Alluvial Architecture, and
1138 the Role of Cohesion in Sinuous River Channels: A Flume Study." *Journal of Sedimentary*
1139 *Research* 77: 197–212.

- 1140 Perkins, R.G., D.M. Paterson, H. Sun, J. Watson, and M.A. Player. 2004. "Extracellular Polymeric
1141 Substances: Quantification and Use in Erosion Experiments." *Continental Shelf Research* 24
1142 (15): 1623–35. doi:10.1016/j.csr.2004.06.001.
- 1143 Puijalon, S., J. Lena, and G. Bornette. 2007. "Interactive Effects of Nutrient and Mechanical
1144 Stresses on Plant Morphology." *Annals of Botany* 100 (6). Oxford University Press: 1297–
1145 1305. doi:10.1093/aob/mcm226.
- 1146 Rajaratnam, N., and K.A. Mazurek. 2002. "Erosion of a Polystyrene Bed by Obliquely Impinging
1147 Circular Turbulent Air Jets." *Journal of Hydraulic Research* 40 (6). Taylor & Francis Group:
1148 709–16. doi:10.1080/00221680209499917.
- 1149 Read, J., and A. Stokes. 2006. "Plant Biomechanics in an Ecological Context." *American Journal of*
1150 *Botany* 93 (10). Botanical Society of America: 1546–65. doi:10.3732/ajb.93.10.1546.
- 1151 Santos, M.G.M., N.P. Mountney, J. Peakall, 2017a "Tectonic and environmental controls on
1152 Palaeozoic fluvial environments: reassessing the impacts of early land plants on
1153 sedimentation." *Journal of the Geological Society, London*, 174, 393–404,
1154 <https://doi.org/10.1144/jgs2016-063>
- 1155 Santos, M.G.M., N.P. Mountney, J. Peakall, R.E. Thomas, P.B. Wignall, D.M. Hodgson, 2017b, "
1156 Reply to Discussion on 'Tectonic and environmental controls on Palaeozoic fluvial
1157 environments: reassessing the impacts of early land plants on sedimentation' *Journal of the*
1158 *Geological Society, London*, <https://doi.org/10.1144/jgs2016-063>
- 1159 Schindler, R.J., D.R. Parsons, L. Ye, J.A. Hope, J.H. Baas, J. Peakall, A.J. Manning, et al. 2015.
1160 "Sticky Stuff: Redefining Bedform Prediction in Modern and Ancient Environments." *Geology*
1161 43 (5). GeoScienceWorld: 399–402. doi:10.1130/G36262.1.
- 1162 Schuurman, F., M.G. Kleinans, and W.A. Marra. 2013. "Physics-Based Modeling of Large
1163 Braided Sand-Bed Rivers: Bar Pattern Formation, Dynamics, and Sensitivity." *Journal of*

- 1164 *Geophysical Research Earth Surface* 118: 2509–2527, doi: 10.1002/2013JF002896.
- 1165 Seal, R, C Paola, G Parker, J B Southard, and P R Wilcock. 1997. “Experiments on Downstream
1166 Fining of Gravel: I. Narrow-Channel Runs.” *J. of Hydraulic Engineering* 123 (10): 874–84.
- 1167 Seminara, G, and M Tubino. 1989. “Alternate Bars and Meandering: Free, Forced and Mixed
1168 Interactions.” In *River Meandering*, edited by S Ikeda and G Parker, 153–80. Washington
1169 D.C.: American Geophysical Union.
- 1170 Sheets, B.A., T.A. Hickson, and C. Paola. 2002. “Assembling the Stratigraphic Record:
1171 Depositional Patterns and Time-Scales in an Experimental Alluvial Basin.” *Basin Research* 14
1172 (3). Blackwell Science Ltd: 287–301, doi:10.1046/j.1365–2117.2002.00185.x.
1173 doi:10.1046/j.1365-2117.2002.00185.x.
- 1174 Shields, A. 1936. “Anwendung der Ähnlichkeitsmechanik und der Turbulenzforschung auf die
1175 Geschiebebewegung.” *Mitteilungen der Preußischen Versuchsanstalt für Wasserbau und*
1176 *Schiffbau*, Berlin, Heft 26.
- 1177 Smith, C E. 1998. “Modeling High Sinuosity Meanders in a Small Flume.” *Geomorphology* 25: 19–
1178 30.
- 1179 Solari, L, and G Parker. 2000. “The Curious Case of Mobility Reversal in Sediment Mixtures.”
1180 *Journal of Hydraulic Engineering* 126 (3): 185–97.
- 1181 Stefanon, L., L. Carniello, A. D’Alpaos, and S. Lanzoni. 2010. “Experimental Analysis of Tidal
1182 Network Growth and Development.” *Continental Shelf Research* 30: 950–62.
- 1183 Struiksmā, N. 1985. “Prediction of 2-D Bed Topography in Rivers.” *Journal of Hydraulic*
1184 *Engineering* 111 (8): 1169–82.
- 1185 Sutherland, J., and R.J.S. Whitehouse. 1998. “Scale Effects in the Physical Modelling of Seabed
1186 Scour.” HR Wallingford.

- 1187 Syvitski, J.P.M., C.J. Vörösmarty, A.J. Kettner, and P. Green. 2005. "Impact of Humans on the
1188 Flux of Terrestrial Sediment to the Global Coastal Ocean." *Science* 308 (5720). American
1189 Association for the Advancement of Science: 376–80. doi:10.1126/science.1109454.
- 1190 Tal, M., K. Gran, A.B. Murray, C. Paola, and D.M. Hicks. 2004. *Riparian Vegetation as a Primary*
1191 *Control on Channel Characteristics in Multi-Thread Rivers*. In: Bennett, S. J. and Simon, A.
1192 *Riparian Vegetation and Fluvial Geomorphology*, Water Sci. Appl. Ser. 8, 43-58, AGU,
1193 Washington, D. C.
- 1194 Tal, M., and C. Paola. 2007. "Dynamic Single-Thread Channels Maintained by the Interaction of
1195 Flow and Vegetation." *Geology* 35 (4). doi:10.1130/G23260A.1.
- 1196 Tal, M., and C. Paola. 2010. "Effects of Vegetation on Channel Morphodynamics: Results and
1197 Insights from Laboratory Experiments." *Earth Surface Processes and Landforms* 35 (9):
1198 1014–28. doi:10.1002/esp.1908.
- 1199 Talmon, A M, N Struiksmā, and M C L M van Mierlo. 1995. "Laboratory Measurements of the
1200 Direction of Sediment Transport on Transverse Alluvial-Bed Slopes." *Journal of Hydraulic*
1201 *Research* 33 (4): 495–517.
- 1202 Tambroni, N., M. Bolla Pittaluga, and G. Seminara. 2005. "Laboratory Observations of the
1203 Morphodynamic Evolution of Tidal Channels and Tidal Inlets." *Journal of Geophysical*
1204 *Research: Earth Surface* 110 (F4). doi:10.1029/2004JF000243.
- 1205 Thomas, R.E., M.F. Johnson, L.E. Frostick, D.R. Parsons, T.J. Bouma, J.T. Dijkstra, O. Eiff, et al.
1206 2014. "Physical Modelling of Water, Fauna and Flora: Knowledge Gaps, Avenues for Future
1207 Research and Infrastructural Needs." *Journal of Hydraulic Research* 52 (3). Taylor & Francis:
1208 311–25. doi:10.1080/00221686.2013.876453.
- 1209 Thorne, C.R., Tovey, N.K., 1981. "Stability of composite river banks" *Earth Surface Processes and*
1210 *Landforms* 6, 469-484

- 1211 Tolhurst, T.J., G. Gust, and D.M. Paterson. 2002. "The Influence of an Extracellular Polymeric
1212 Substance (EPS) on Cohesive Sediment Stability." In *Proceedings in Marine Science*, 5:409–
1213 25. doi:10.1016/S1568-2692(02)80030-4.
- 1214 Toro-Escobar, C, C Paola, G Parker, P R Wilcock, and J B Southard. 2000. "Experiments on
1215 Downstream Fining of Gravel. II: Wide and Sandy Runs." *J. of Hydraulic Engineering* 126
1216 (3): 198–208.
- 1217 Tsujimoto, T. 1990. "Distorted Model and Time Scale Evaluation of Multiscale Subjected Fluvial
1218 Processes." In *Movable Bed Physical Models*, 31–48. Springer.
- 1219 Turowski, J.M., D. Lague, A. Crave, and N. Hovius. 2006. "Experimental Channel Response to
1220 Tectonic Uplift." *Journal of Geophysical Research: Earth Surface* 111 (F3). Wiley Online
1221 Library.
- 1222 Van de Lageweg, W.I., W.M. Van Dijk, D. Box, and M.G. Kleinhans. 2016. "Archimetrics: A
1223 Quantitative Tool to Predict Three-Dimensional Meander Belt Sandbody Heterogeneity." *The*
1224 *Depositional Record* 2 (1): 22–46. doi:10.1002/dep2.12.
- 1225 Van de Lageweg, W I, W M Van Dijk, R Hoendervoogt, and M G Kleinhans. 2010. "Effects of
1226 Riparian Vegetation on Experimental Channel Dynamics." In *Riverflow 2010, Volume 2*,
1227 edited by Dittrich, Koll, Aberle, and Geisenhainer, 1331–38. Bundesanstalt für Wasserbau.
- 1228 Van de Lageweg, W.I., S.J. Mclelland, D.R. Parsons, 2017, "Quantifying biostabilisation effects of
1229 biofilm-secreted and synthetic extracellular polymeric substances (EPS) on sandy substrate"
1230 *Earth Surface Dynamics Discussions*
- 1231 Van Dijk, W.M., W.I. Van de Lageweg, and M.G. Kleinhans. 2012. "Experimental Meandering
1232 River with Chute Cutoffs." *Journal of Geophysical Research Earth Surface* 117: F03023,
1233 doi:10.1029/2011JF002314.
- 1234 Van Dijk, W.M., R. Teske, W.I. van de Lageweg, and M.G. Kleinhans. 2013a. "Effects of

- 1235 Vegetation Distribution on Experimental River Channel Dynamics.” *Water Resources*
1236 *Research* 49 (11): 7558–7574, doi:10.1002/2013WR013574. doi:10.1002/2013WR013574.
- 1237 Van Dijk, W.M., W.I. Van de Lageweg, and M.G. Kleinhans. 2013b. “Formation of a Cohesive
1238 Floodplain in a Dynamic Experimental Meandering River.” *Earth Surface Processes and*
1239 *Landforms* 38: 1550–1565, doi:10.1002/esp.3400.
- 1240 van Oorschot, M., M.G. Kleinhans, G. Geerling, and H. Middelkoop. 2016. “Distinct Patterns of
1241 Interaction between Vegetation and Morphodynamics.” *Earth Surface Processes and*
1242 *Landforms* 41 (6): 791–808. doi:10.1002/esp.3864.
- 1243 Van Rijn, L. C., P. K. Tonnon, A. Sanchez-Arcilla, I. Cáceres, and J. Grüne. 2011. “Scaling Laws
1244 for Beach and Dune Erosion Processes.” *Coastal Engineering* 58 (7): 623–36.
1245 doi:10.1016/j.coastaleng.2011.01.008.
- 1246 Vollmer, S., and M.G. Kleinhans. 2007. “Predicting Incipient Motion, Including the Effect of
1247 Turbulent Pressure Fluctuations in the Bed.” *Water Resources Research* 43 (5): W05410,
1248 doi:10.1029/2006WR004919. doi:10.1029/2006WR004919.
- 1249 Vollmers, H., and E. Giese. 1972. “Elbe Tidal Model with Movable Bed.” In *Coastal Engineering*
1250 *1972*, 2457–73. New York, NY, NY: American Society of Civil Engineers.
1251 doi:10.1061/9780872620490.144.
- 1252 Wang, X.L., F.W. Zwiars, and V.R. Swail. 2004. “North Atlantic Ocean Wave Climate Change
1253 Scenarios for the Twenty-First Century.” *Journal of Climate* 17 (12): 2368–83.
1254 doi:10.1175/1520-0442(2004)017<2368:NAOWCC>2.0.CO;2.
- 1255 Wilcock, P.R., C.H. Orr, and J.D.G. Marr. 2008. “The Need for Full-Scale Experiments in River
1256 Science.” *Eos, Transactions American Geophysical Union* 89 (1): 6–6.
1257 doi:10.1029/2008EO010003.
- 1258 Wilcock, P R. 1993. “Critical Shear Stress of Natural Sediments.” *Journal of Hydraulic*

- 1259 *Engineering* 119 (4): 491–505.
- 1260 Willson, C.I.S., N. Dill, W. Barlett, S. Danchuk, and R. Waldron. 2007. “Physical and Numerical
1261 Modeling of River and Sediment Diversions in the Lower Mississippi River Delta.” In *Coastal*
1262 *Sediments '07*, 749–61. Reston, VA: American Society of Civil Engineers.
1263 doi:10.1061/40926(239)57.
- 1264 Wobus, C.W., G.E. Tucker, and R.S. Anderson. 2010. “Does Climate Change Create Distinctive
1265 Patterns of Landscape Incision?” *Journal of Geophysical Research: Earth Surface* 115 (F4).
1266 Wiley Online Library.
- 1267 Yager, E.M., M. Kenworthy, and A. Monsalve. 2015. “Taking the River inside: Fundamental
1268 Advances from Laboratory Experiments in Measuring and Understanding Bedload Transport
1269 Processes.” *Geomorphology* 244 (September): 21–32. doi:10.1016/j.geomorph.2015.04.002.
- 1270 Yalin, M.S. 1959. “Über die Naturähnlichkeit der Geschiebebewegung bei Modellversuchen.” *Die*
1271 *Bautechnik*, 36. Jahrgang, no. 3: 96–99.
- 1272 Yalin, M S. 1971. *Theory of Hydraulic Models*. London, UK: Macmillan.
- 1273 Yalin, M S, and J W Kamphuis. 1971. “Theory of Dimensions and Spurious Correlation.” *Journal*
1274 *of Hydraulic Research* 9 (2). Taylor & Francis Group: 249–65.
- 1275 Young, W J, and T Davies. 1991. “Bedload Transport Processes in a Braided Gravel-Bed River
1276 Model.” *Earth Surface Processes and Landforms* 16: 499–511.
- 1277 Zanke, U.C.E. 2003. “On the Influence of Turbulence on the Initiation of Sediment Motion.”
1278 *International Journal of Sediment Research* 18 (1): 1–15.
- 1279 Zhao, G., P.J. Visser, J. Lu, and J.K. Vrijling. 2013. “Similarity of the Velocity Profile in
1280 Geometrically Distorted Flow Model.” *Flow Measurement and Instrumentation* 32 (August).
1281 Elsevier: 107–10. doi:10.1016/J.FLOWMEASINST.2013.04.005.

- 1282 Zwamborn, J.A. 1966. "Reproducibility in Hydraulic Models of Prototype River Morphology." *La*
1283 *Houille Blanche*, no. 3(April). EDP Sciences: 291–98. doi:10.1051/lhb/1966020.
1284