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Probing the Molecular Orientation of a Single Conjugated Polymer via Nano-gap SERS

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ABSTRACT: Determining the molecular orientation at the single molecule level is of key importance for a wide range of applications ranging from molecular electronic devices to biomedical applications. In this work Surface-Enhanced Raman Scattering (SERS) was used to probe the light emitting conjugated polymer F8-PFB at the single molecule level using nanoparticle on an extended metallic film nano-gap. The directional field enhancement of the nano-gap combined with Density Functional Theory (DFT) calculations was used to determine the orientation of the molecule. This analysis revealed that the spin-coated conjugated polymer preferentially aligns its molecular chains parallel to the metallic substrate. The integration of this approach in nanofabrication and synthesis will have a profound impact on different fields ranging from molecular electronic devices to biomedical applications.

INTRODUCTION

Conjugated polymers have been of significant interest due to their applications in optoelectronic devices such as organic light emitting diodes, photovoltaics, lasers, as well as bio-probes and sensors.¹⁻³ The optical and electronic properties of conjugated polymer thin-films are highly dependent on the processing and deposition conditions.^{4,5} illustrating the strong relationship between morphology and photophysics. This relationship can be revealed by studying these materials at the single molecule level which is commonly done using single molecule fluorescence spectroscopy.⁶ Although this technique provides a wealth of information regarding the processes that contribute to excited state relaxation and energy transfer between chromophores of large molecules, it brings limited structural/conformational information about these materials. Raman spectroscopy, on the other hand, has been shown to be a





powerful technique for materials characterization, providing detailed structural information in a label free and noninvasive manner.^{7,8}

The use of Raman scattering for molecular identification down to the single molecule level can be challenging and requires large enhancements of the electromagnetic (EM) fields in the analyte local environment. By coupling Raman spectroscopy with plasmonic devices, SERS (Surface Enhanced Raman Scattering) has been shown to provide molecular fingerprints unique to the material under investigation down to the single molecule regime.⁹⁻¹³

On the other hand, determining the molecular orientation of an analyte *in situ* is of key importance in a range of applications. In optoelectronics, molecular orientation is important for the development of high performance devices, in particular of organic light-emitting diodes (OLEDs), organic field-effect transistors (OFETs) and molecular switches.^{4,14,15} Furthermore, in biomedical sciences, knowing the molecular orientation enables conformation changes in DNA to be monitored,¹⁶ as well as probing in more details the forward rates of cell-bound receptor-ligand binding which has been shown to be strongly influenced by molecular orientation.^{17,18} Additionally, by understanding the interaction mechanisms of biomolecules the possibility for drug development also increases significantly.¹⁹

Polyfluorene (PF) based materials generated special attention, and the importance of the chain orientations with respect to the substrate were highlighted by various studies. ^{20,21} In bulk films this orientation can be recovered by specular or grazing incidence x-ray diffraction for polycrystalline samples,²² or by polarized techniques (optical absorption, photoluminescence, Raman or x-ray microscopy).^{21,23}

There is however an increasing need to probe molecular orientations in materials at a single molecule level. This can be done by azimuthally or radially polarized laser beams,^{24,25} or by SERS. The orientation of small molecular targets through SERS has already been recovered in previous studies.^{19,26-32} However, using SERS to determine the molecular orientation of large molecules such as conjugated polymers is still a real challenge as for large molecules the molecular sizes are typically larger than the enhancement volume. Additionally molecular diffusion leads to changes in molecular conformations and orientations within the enhancement volume resulting in selective Raman enhancement over time. These two effects combine to make determining the orientation of large molecules more challenging.^{33,34}

Up to now only two studies report on SERS spectra of single conjugated polymer chains,^{35,36} where simple nanostructured silver films were used to provide the necessary electromagnetic enhancement. However, the random nature of the enhanced electromagnetic field associated with these films averages out the inherited information on molecular orientation contained within the SERS signal. The nano-gap architecture used throughout this study (see Figure 1 and Figure S1)) was shown to preferentially enhance Raman modes which are aligned with the nano-gap field.³² This preferential alignment in the enhanced Raman modes allows to determine the orientation of the molecule relative to the nano-gap field.

In this work, the orientation of a single chain of the PFbased conjugated co-polymer, poly(9,9'-dioctylfluorene-cobis-N,N'-(4-butylphenyl)-bis-N,N'-phenyl-1,4-phenylenediamine), otherwise known as F8-PFB (0.95:0.05 F8:PFB and see Figure 2(a)) was recovered using the well-defined electromagnetic field enhancement direction in a plasmonic nano-gap formed between a silver nanoparticles and a silver film (see Figure 1). Silver nanoparticles of 118 nm in diameter were coupled to a 100 nm thick silver film to create the plasmonic nano-gaps which have been shown to be optimized for SERS excitation at a wavelength of 532nm.³² By combining the single molecule SERS spectra from those nano-gaps with Density Functional Theory (DFT) calculations, the molecular orientation of a single chain of F8:PFB was determined. The results of this analysis clearly demonstrate that spin-coated conjugated polymers tend to align their molecular chains parallel to the metallic substrate.

METHODS

The plasmonic nano-gaps were fabricated on glass slides which were first sequentially cleaned in solutions of acetone, isopropanol and deionized water for 10 minute periods in an ultrasonic bath. After cleaning, the slides were thermally coated with a silver film of average thickness 100 nm, forming the bottom of the nano-gap (Figure 1). The conjugated polymer F8-PFB ($M_w = 355$ kDa and 0.95:0.05 F8:PFB), dissolved in toluene at concentrations ranging between 2.8×10⁻⁵ M to 2.8×10⁻¹⁰ M (ref. 6,13,35), was deposited via spin coating. Lastly, silver nanoparticles of average diameter 118 nm (Sigma-Aldrich) were spin-coated onto the surface of the polymer coated film from a solution of ethanol at a concentration of 2.7×10⁶ particle.L⁻¹. This concentration was used to obtain several microns spacing between the nanoparticles (see Figure S2), therefore ensuring that the nano-gaps are isolated from one another.

SERS measurements were performed using a home built Raman setup with a 532 nm laser line using off-resonance excitation (see Figure 2(b)). The measurements were acquired via direct excitation through a 50× magnification objective lens with Numerical Aperture NA = 0.55 (Mitutoyo M Plan Apo 50X). The collected Raman signal was then focused onto the entrance slit of an iHR320 Horiba spectrometer where it was dispersed using a 1200 line.mm⁻¹ grating onto a liquid nitrogen cooled Symphony CCD. Rayleigh scattering was filtered from the Stokes scattered light using two Kaiser holographic notch and super-notch filters centered at a wavelength of 532 nm.

The experimental Raman scattering spectra were complemented with DFT calculations, done on $F8_n$ oligomers (n = 1,2,3,6), PFB and the F8₃PFB co-oligomer. The geometries of F8_n were optimized in the C₂ point group. For all calculations, the B3LYP functional and 6-311G* basis set were employed as defined in the Gaussian 09 program package.³⁷ As the PFB unit contributes only very little to the Raman spectrum (vide infra), and in experiment the PFB content is only 5%, interpretation of the Raman spectra were based on the $F8_n$ oligomers. Raman frequencies v_i at the $F8_\infty$ polymer limit were obtained by an exponential extrapolation of v_i vs. N⁻¹, i.e. the inverse number of double bonds along the shortest path between the terminal C-atoms.³⁸ As the frequencies of the relevant modes at the polymer limit are practically identical with those of F86 (vide infra), the analysis of the Raman tensors relied on F86. For this, the absolute values of the atomic displacements (see Figure S2) were summed, weighted by the square root of the mass; using this information, the inclination against the axis defined relative to the terminal carbon atoms in the chain was resolved.

RESULTS AND DISCUSSION

Figure 2(c) shows the Raman signal from F8-PFB at a concentration of 2.8×10^{-5} M on a glass substrate which provides a full fingerprint spectrum of the material. The most prominent Raman mode of the material is observed at 1605 cm⁻¹ which, using DFT calculations on F8₆, can be attributed to a mode corresponding to combined phenyl inter- and intraring C–C stretches (*vide infra*). Additional, albeit less intense Raman modes can be identified in the range between 1100-1500 cm⁻¹ (with prominent bands at 1136, 1258, 1281, 1350, 1421 cm⁻¹). The DFT analysis assigns these bands to normal modes combining intra- and inter-ring C-C stretching, and C-H and C-C bending motions. than 3x10⁻⁹ M, the intensity of the Raman signal no longer decreases with concentration but remains constant (Figure 3(b)), suggesting that, for such low concentrations, the Raman signal is attributed to a single (macro)molecule in the nano-gap.



Figure 2. (a) The chemical structure of F8-PFB. (b) F8-PFB Absorption (black line) and Emission (red dotted line) spectra peaking at 383 nm and 463 nm respectively. (c) F8-PFB Raman spectrum taken from a thin film on glass at 2.8×10^{-5} M.

The SERS spectrum of F8-PFB was then studied inside the nano-gaps at different concentrations. As the concentration of F8-PFB is systematically lowered with each sample, a corresponding decrease in intensity of the SERS peaks is observed and attributed to the reducing number of molecules contributing to the Raman signal (Figure 3(a)). The decrease in concentration also reduces the film thickness (Figure 3(a)) which we measured at each concentration using surface profilometry. However, for concentrations lower



Figure 3. SERS concentration study of F8-PFB for 118 nm silver nanoparticles on a silver film (estimated thickness in parentheses) taken at 532 nm excitation with 3 s collection time. (a) The full concentration range and (b), the point at which intensity no longer drops.

To further confirm the single molecule regime, 10 consecutive SERS measurements were taken from nano-gaps with two different polymer concentrations, 2.8×10^{-8} M and 2.8×10^{-9} M, over a period of 30 seconds. As can be seen in Figure 4, for the higher concentration of 2.8×10^{-8} M, the intensity of the 1605cm⁻¹ C–C stretching mode is constant. However, lowering the concentration one order of magnitude to 2.8×10^{-9} M, leads to the Raman signal to blink in intensity over time. In the configuration considered here, with the molecule in direct contact with metallic surfaces, this characteristic behavior of single molecule events can be explained by photoionization via charge-transfer states.³⁹⁻⁴¹



Figure 4. 10 consecutive SERS measurements for F8-PFB at (a) 2.8×10^{-8} M and (b) 2.8×10^{-9} M. The C-C stretching mode fluctuates in intensity with each successive measurement.

To explore the single molecule Raman spectra of F8-PFB further, typical single molecule SERS spectra of F8-PFB measured from nano-gaps at concentration 2.8x10⁻⁹ M are shown in Figure 5 (panels 2-6), along with the corresponding Raman spectrum at high concentration for comparison (panel 1). In contrast to the high concentration Raman spectrum, which shows one strong peak at 1605 cm⁻¹, the single molecule data clearly shows two Raman peaks, centered at 1605 cm⁻¹ and 1556 cm⁻¹ respectively.

The DFT calculations demonstrate that the two peaks observed experimentally are dominated each by single vibrations (see Figure 6(a)); furthermore, PFB vibrations are of low intensity (see inset of Figure 6(a)) and can be thus safely neglected. The R₁ frequency remains fairly constant with n, giving 1617 cm⁻¹ at the polymer limit (see Figure 6(b)); it combines intra-fluorene stretching and asymmetric breathing motions (Figure 7(a)) and is polarized at a 46° angle to the terminal C-atoms of the polyfluorene chain (see Figure 7(b)). R₂ decreases with n, but saturates fast towards the polymer limit, being essentially the same as for n = 6 (1642 cm⁻¹). R₂ combines fluorene inter- and intra-ring C-C stretches (Figure 7(a)) and is polarized at a 17° angle to the terminal C-atoms of the polyfluorene chain (Figure 7(b)); the enclosed angle between R_1 and R_2 is 50° (Figure 7(c)).



Figure 5. Raman spectra from F8-PFB at 2.8×10^{-5} M on glass (panel 1), and SERS spectra from F8-PFB at 2.8×10^{-9} M within the nano-gaps between a silver nanoparticle and silver film (panels 2-6).

As mentioned earlier, in SERS, the relative intensities of the peaks depend on the molecular orientation relative to the nano-gap field. For each analyte molecule, the SERS signal intensity, I_{SERS} is proportional to $\cos^2(\theta)$,⁴² and can be written as follows:

$$I_{SERS} = I_{Raman} \cdot G \cdot \cos^2 \theta \ (1)$$

Where θ is the angle between the driving electric field and the direction of Raman tensor of the mode (Figure 7b), I_{Raman} is the intensity of the Raman mode in the absence of plasmonic structure, and *G* is the SERS enhancement factor.³²

Using equation (1), the SERS enhancements in R_1 $I_{SERS_{R_1}}$ and R_2 $I_{SERS_{R_2}}$ can be written as:

$$I_{SERS_{R1}} = I_{DFT_{R1}} \cdot G \cdot \cos^2(\theta) (2)$$
$$I_{SERS_{R2}} = I_{DFT_{R2}} \cdot G \cdot \cos^2(\theta + 50)(3)$$

The ratio between $I_{SERS_{R1}}$ and $I_{SERS_{R2}}$ can therefore be expressed as:

$$\frac{I_{SERS_{R1}}}{I_{SERS_{R2}}} = \frac{I_{DFT_{R1}}}{I_{DFT_{R2}}} \cdot \frac{\cos^2(\theta)}{\cos^2(\theta + 50)}$$
(4)



Figure 6. DFT calculations demonstrating (a) Raman spectrum of F8₆ (PFB as inset). (b) Peak shift of R₁ and R₂ Raman frequencies towards the polymer limit ($N^{-1} = 0$; see Method section).

By combining equation (4) with the single molecule spectra of figure 5, it is possible to determine the angle θ , and therefore the orientation of the single polymer chain within each nano-gap (table 1).

Table 1. Calculating the orientation of the target analyte with respect to the cavity using experimental SERS intensity for peaks R_1 and R_2 at single molecule spectra from Figure 5.

Intensity R ₁	Intensity R ₂	$\theta(R_1)$	θ(R ₂)
328	337	28°	78°
589	315	35°	85°
207	305	23°	73°
256	330	25°	75°
342	234	32°	82°

Table 1 shows that the angle for the R_2 band calculated from the DFT data and SERS spectra is between 73° to 85° relative to nano-gap field (z-axis in figure 7c). On average R_2 makes an angle of 11° relative to the substrate. Therefore, when considering the main axis of the R_2 tensor in figure 7b as being 17° with respect to the long polymer axis, our calculations indicate that the spin-coated F8-PFB molecules align their molecular chain almost parallel to the substrate within the nano-gap. This is further supported by the comparable intensities of R_1 and R_2 in SERS which arise from their respective projection of their dipole derivative vectors along the nano-gap field.



Figure 7. (a) Relative atomic displacements of the Catoms in the central fluorene units for R_1 and R_2 . (b) Main axis of the polarizability tensor of R_1 and R_2 for F8₆: (angles vs. terminal C-C). (c) Orientation of the target analyte with respect to the nano-gap field.

It should be stressed that the high in-plane orientation found in the current work is much larger than in ellipsometry studies e.g. on spin-coated poly-dicoctylfluorene films.²¹ We attribute this mainly to the much thicker films in these experiments; furthermore the substrate might have a significant effect. The observed differences high-light the advantage of the SERS-based method; i.e. to obtain local information on polymer organization at a single molecular level.

CONCLUSION

In conclusion, we have shown that through combining SM-SERS spectroscopy and DFT calculations, it is possible to determine the molecular orientation of conjugated polymers from plasmonic nano-gaps at concentrations demonstrating single molecule level characteristics. Our analysis revealed that spin-coated conjugated polymers preferentially align their molecular chains almost parallel to the substrate. The work shown in this study provides a better understanding of the interfacial interaction of conjugated polymers with metallic surfaces and offers a powerful tool to optimizing conjugated polymers based devices such as OLEDs, OFETs and molecular switches. In parallel, *in situ* retrieval of the molecular orientation and confirmation will have profound impacts on biomedical sciences, allowing for, amongst others, DNA monitoring, probing of forward rates of cell-bound receptor-ligand binding, and enhanced drug development and delivery.

ASSOCIATED CONTENT

Supporting Information.

The Supporting Information is available free of charge via the Internet at <u>http://pubs.acs.org</u>."

Calculated field enhancement, typical dark-field imaging of the sample, calculated Raman modes for F8-PFB (pdf)

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Author Contributions

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Notes

The Authors declare no competing financial interest.

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ABBREVIATIONS

SM; Single Molecule; SERS, Surface Enhanced Raman Spectroscopy; DFT, Density functional Theory.

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